

Information Sciences Letters An International Journal

http://dx.doi.org/10.18576/isl/090305

An efficient approximate-analytical method to solve time-fractional KdV and KdVB equations

Hegagi Mohamed Ali^{1,*}

¹Department of Mathematics, Faculty of Science, Aswan University, Aswan 81528, Egypt

Received: 3 May 2020, Revised: 9 July 2020, Accepted: 13 July 2020

Published online: 1 Sep. 2020

Abstract: In this article, we present the modified generalized Mittag-Leffler function method (MGMLFM) as an approximate-analytical method to give a proper solution of time-fractional Korteweg-de Vries (KdV) and Korteweg-de Vries-Burger's (KdVB) equations, which have various applications in physics and applied mathematics. The time-fractional partial derivatives are based on Caputo sense. The obtained solution is constructed in a rapidly convergent power series. By comparing the approximate MGMLFM solutions when the fractional operator equal one with known exact solutions we have an appropriate agreement. The advantage of the article is to apply the suggested method to solve linear and nonlinear time-fractional partial differential equations, where it needs less computational effort which saves time and effort. The convergence of absolute error be controlled on by the parameters in the time-fractional KdV and KdVB equations were found. The simulation of the obtained results is presented in the forms of graphs to illustrate the reliability and efficiency of our method.

Keywords: Applications in the physical sciences; Time-fractional partial differential equations; Mittag-Leffler function; Series solutions; KdV and KdVB equations.

1 Introduction

Korteweg de Vries (KdV) equation was first proposed by Korteweg and de Vries [1]. After that, the KdV equation has been contributing to describe various nonlinear phenomena in physics and applied mathematics such as solid-state physics, particle acoustic waves, stratified internal waves, plasma physics, particle acoustic waves, fluid mechanics, quantum field and so on [2,3,4,5]. The standard form of the KdV equation is

$$U_t + \lambda U U_x + \mu U_{xxx} = 0, \tag{1}$$

where λ and μ are real constants not equal to zero.

The Korteweg de Vries-Burgers (KdVB) equation was introduced by Su and Gardner [6], which considered as a combination of the KdV equation (1) and the following Burgers equation

$$U_t + \lambda U U_x + \epsilon U_{xx} = 0, \tag{2}$$

where ϵ is real constant not equal to zero. So, the standard form of the KdVB equation is

$$U_t + \lambda U U_x + \epsilon U_{xx} + \mu U_{xxx} = 0. \tag{3}$$

The KdVB equation has been used as a nonlinear model to describe physical phenomena of interest such as plasma waves, the flow of liquids containing gas bubbles, the propagation of waves on an elastic tube filled with a viscous fluid, the propagation of undular bores in shallow water, and turbulence, etc (see e.g., [7,8,9,10,11,12,13,14]).

There are many methods used to provide numerical and analytical solutions for KdV and KdVB equations in the literature such as decomposition and Adomian decomposition method [15,16], Radial basis function [17, 18], tanh method [19,20], Homotopy and Homotopy perturbation method [21,22], Finite difference method [23,24], Ansatz and cubic B-spline Galerkin method [25, 26], and for other different methods (see e.g. [27,28,29, 30,31]).

Recently, fractional partial differential equations (FPDEs) have appeared in numerous research fields of physics, finance, engineering and many other applied science (see e.g., [32,33,34,35,36,37,38,39,40,41,42]) and some references cited therein. This is due to the accuracy and realistic of FPDEs than others in describing mathematical models that describe these phenomena in

^{*} Corresponding author e-mail: hegagi_math@aswu.edu.eg



these various fields, where depend on upon all the historical states not only its current state. There are many researchers have devoted their efforts to develop methods for solving linear and nonlinear FPDEs and it is noted that there is difficulty in finding approximate and analytical methods of solution, therefore, to find accurate methods to solve these FPDEs are still under consideration. For instance, some of the methods solving FPDEs as homotopy analysis method (HAM) [43,44,45, 46,47], homotopy perturbation method (HPM)[48,49,50, 51], generalized Mittag-Leffler function (GMLFM) [54,55,56,57,58], homotopy transform method (HATM) [52,53] and so on.

The main goal and motivation of this research is to investigate MGMLFM as an approximate-analytical method to solve the time-fractional KdV and KdVB equations in the following form, respectively,

$${}_{0}^{C}D_{t}^{\alpha}U(x,t) + U(x,t)U_{x}(x,t) + \frac{1}{2}U_{xxx}(x,t) = 0,$$
 (4)

and

$${}_{0}^{C}D_{t}^{\alpha}U(x,t) + U(x,t)U_{x}(x,t) - \delta U_{xx}(x,t) + \frac{1}{2}U_{xxx}(x,t) = 0,$$
(5)

where ${}_{0}^{C}D_{t}^{\alpha}$ is a Caputo fractional partial derivative of order α with respect to time, δ is real constant not equal to zero, U_{xx} represents a viscous loss, U_{xxx} is dispersion and UU_x is convective nonlinearity.

The most advantages of this method is an easy and simple technique to solve linear and nonlinear FPDEs and achieves the work with small efforts of computational comparing with other methods. We compare the solutions obtained by MGMLFM with the exact solutions and extract the absolute error to prove the efficiency of the method. Moreover, Eqs. (4) and (5) have never been studied before by the proposed method which indicates the novelty of results.

This paper is organized as follows. In Section 2, we give some basic definitions and preliminaries facts which are required to reach our main results. In Section 3, we introduce the analysis of the MGMLFM to solve a general form of time-FPDEs. Section 4 is divided into two subsections and its devoted to applying the MGMLFM to solve time-fractional KdV and KdVB equations. In Section 5, we offered some numerical simulations to compare our results with the exact solution in order to prove the accuracy and efficacy of our methodology. Finally, Section 6 presents conclusion of this research.

2 Some Basic Definitions and Preliminaries **Facts**

This section presents a review of some basic concepts which are essentially relevant to the results of this manuscript (see e.g. [59,60]).

Definition 1.Let $\mathcal{F}(t)$ be an integrable function on the interval $[t_0,T]$, $t \in [t_0,T]$. Then, the Riemann-Liouville fractional integral of order $\alpha > 0$ defined by

$$t_0 I_t^{\alpha} \mathcal{F}(t) = \frac{1}{\Gamma(\alpha)} \int_{t_0}^t (t - \eta)^{\alpha - 1} \mathcal{F}(\eta) d\eta, \quad t_0 \ge 0, \quad t > t_0,$$
$$t_0 I_t^0 \mathcal{F}(t) = \mathcal{F}(t),$$

where $\Gamma(\cdot)$ is the Euler gamma function which defined as follows

$$\Gamma(\xi) = \int_0^\infty t^{\xi - 1} e^{-t} dt, \quad (Re(\xi) > 0).$$

Definition 2.Let $\mathcal{F}(x,t)$ be absolutely continuous functions on the interval $[t_0,T],\ t\in [t_0,T].$ and $n < \alpha \leq n-1$, where $n \in \mathbb{N}$. Then, the Caputo fractional partial derivative is defined as

$${}_{t_0}^C D_t^{\alpha} \mathcal{F}(x,t) = \frac{1}{\Gamma(n-\alpha)} \int_{t_0}^t (t-\eta)^{n-\alpha-1} \frac{\partial^n \mathcal{F}(x,\eta)}{\partial \eta^n} d\eta,$$
$$t_0 \ge 0, \ t > t_0,$$

In particularly, for $0 < \alpha < 1$, the Caputo fractional partial derivative becomes

$${}_{t_0}^C D_t^{\alpha} \mathcal{F}(x,t) = \frac{1}{\Gamma(1-\alpha)} \int_{t_0}^t (t-\eta)^{-\alpha} \frac{\partial \mathcal{F}(x,\eta)}{\partial \eta} d\eta,$$
$$t_0 \ge 0, \quad t > t_0.$$

Theorem 1.let $\mathcal{F}(x,t)$ be a differentiable function in the interval $[t_0, T]$, $n-1 < \alpha \le n$, $n \in \mathbb{N}$, and $\gamma > -1$, then

$$_{t_0}^C D_{t_0}^{\alpha} I_{t_0}^{\alpha} \mathcal{F}(x,t) = \mathcal{F}(x,t),$$

$${}_{t_0}I_t^\alpha {}_{t_0}^C D_t^\alpha \mathcal{F}(x,t) = \mathcal{F}(x,t) - \sum_{k=0}^{n-1} \frac{\partial^k \mathcal{F}(x,t_0)}{\partial t^k} \frac{(t-t_0)^k}{k!}.$$

In addition, the fractional operator satisfies the following properties:

$${}_{0}^{\alpha}D_{t}^{\alpha}t^{\gamma} = \frac{\Gamma(\gamma+1)}{\Gamma(\gamma-\alpha+1)}t^{\gamma-\alpha},$$

$${}_{0}I_{t}^{\alpha}t^{\gamma} = \frac{\Gamma(\gamma+1)}{\Gamma(\gamma+\alpha+1)}t^{\gamma+\alpha}.$$

Definition 3.The two-parameter Mittag-Leffler function defined by the power series in the form

$$E_{\alpha,\beta}(x) = \sum_{n=0}^{\infty} \frac{x^n}{\Gamma(n\alpha + \beta)}, \quad \alpha, \beta > 0,$$

if $\beta = 1$, this function is denoted by $E_{\alpha}(\cdot)$, and if $\alpha = \beta =$ 1 this function represent the exponential function.

Lemma 1.The Caputo fractional of generalized Mittag-Leffler function is given by



$${}^{C}_{0}D^{\alpha}_{t}E_{\alpha}(\lambda t^{\alpha}) = {}^{C}_{0}D^{\alpha}_{t}(\sum_{n=0}^{\infty} \frac{\lambda^{n}t^{n\alpha}}{\Gamma(n\alpha+1)}) = \sum_{n=1}^{\infty} \frac{\lambda^{n}t^{(n-1)\alpha}}{\Gamma((n-1)\alpha+1)} \qquad \mathcal{N}(u(x,t)) = \mathcal{N}(\sum_{j=0}^{\infty} \mathcal{G}(x)A^{j}\frac{t^{j\alpha}}{\Gamma(j\alpha+1)})$$

$$= \sum_{n=0}^{\infty} \frac{\lambda^{n+1}t^{n\alpha}}{\Gamma(n\alpha+1)} = \lambda E_{\alpha}(\lambda t^{\alpha}).$$

$$= \mathcal{N}(\sum_{j=0}^{\infty} \mathcal{G}(x)u_{j}(t))$$

Theorem 2.[61] Let a function $u(x,t) = \sum_{k=0}^{\infty} \xi^k u_k(x,t)$, then a nonlinear operator $\mathcal{N}(u)$ satisfies the following

$$\frac{\partial^n}{\partial \xi^n} \mathcal{N}(u)_{\xi=0} = \frac{\partial^n}{\partial \xi^n} \mathcal{N}\left(\sum_{k=0}^\infty \xi^k u_k\right)_{\xi=0} = \frac{\partial^n}{\partial \xi^n} \mathcal{N}\left(\sum_{k=0}^n \xi^k u_k\right)_{\xi=0}$$

3 Description of the MGMLFM

In this section, we demonstrate the fundamental idea of MGMLFM to solve nonlinear FPDEs with initial conditions of the following general form:

$${}_{to}^{C}D_{t}^{\alpha}u(x,t) = \mathcal{L}(u(x,t)) + \mathcal{N}(u(x,t)), \qquad (6)$$

with the initial condition

$$u(x,0) = \mathcal{G}(x),\tag{7}$$

where \mathcal{L} and \mathcal{N} are constituted the general linear and nonlinear differential operator for the function u(x,t), respectively, and $\mathcal{G}(x)$ is a know function of variable x.

The MGMLFM suggested that the solution of Eq.(6) can be written as the following

$$u(x,t) = G(x)E_{\alpha}(At^{\alpha}) = \sum_{j=0}^{\infty} G(x)A^{j} \frac{t^{j\alpha}}{\Gamma(j\alpha+1)}, (8)$$

where A is undetermined coefficient, from the initial condition (7), we obtain $G(x) = \mathcal{G}(x)$. Furthermore, by using Lemma 1 the nonlinear FPDE (6) becomes

$$\begin{split} \sum_{j=0}^{\infty} \mathcal{G}(x) A^{j+1} \frac{t^{j\alpha}}{\Gamma(j\alpha+1)} &= L(\sum_{j=0}^{\infty} \mathcal{G}(x) A^{j} \frac{t^{j\alpha}}{\Gamma(j\alpha+1)}) \\ &+ N(\sum_{j=0}^{\infty} \mathcal{G}(x) A^{j} \frac{t^{j\alpha}}{\Gamma(j\alpha+1)}) (9) \end{split}$$

Therefore, the linear part can be decomposed as

$$\mathcal{L}(u(x,t)) = \mathcal{L}(\sum_{j=0}^{\infty} \mathcal{G}(x) A^{j} \frac{t^{j\alpha}}{\Gamma(j\alpha+1)})$$

$$= \mathcal{L}(\mathcal{G}(x)) \sum_{j=0}^{\infty} A^{j} \frac{t^{j\alpha}}{\Gamma(j\alpha+1)}$$

$$= h \mathcal{G}(x) \sum_{j=0}^{\infty} A^{j} \frac{t^{j\alpha}}{\Gamma(j\alpha+1)}, \qquad (10)$$

where h is a constant. By helping of the Theorem 2 and He's polynomials [61,62,63] the nonlinear part can be written as follows

$$\mathcal{N}(u(x,t)) = \mathcal{N}\left(\sum_{j=0}^{\infty} \mathcal{G}(x)A^{j} \frac{t^{j\alpha}}{\Gamma(j\alpha+1)}\right)$$

$$= \mathcal{N}\left(\sum_{j=0}^{\infty} \mathcal{G}(x)u_{j}(t)\right)$$

$$= \mathcal{N}\left(\mathcal{G}(x)\right)\left(\mathcal{N}(u_{0}(t)) + \sum_{j=1}^{\infty} \left(\mathcal{N}\left(\sum_{k=0}^{j} u_{k}(t)\right)\right)\right)$$

$$-\mathcal{N}\left(\sum_{k=0}^{j-1} u_{k}(t)\right)\right)$$
(11)

By replacing linear and nonlinear parts from Eq.(10) and Eq.(11) in Eq.(9), this leading to identify the recurrence relation and obtain the coefficient A and then obtained the general solution of the nonlinear FPDE.

4 Implementation of the MGMLFM and Results

Here, we apply the MGMLFM on KdV and KdVB equations involving on the time-fractional derivative. Furthermore, we present some numerical simulations to illustrate the advantages and accuracy of the proposed method.

4.1 Time-Fractional KdV Equation

We consider the time-fractional KdV equation as follows

$${}_{0}^{C}D_{t}^{\alpha}U(x,t) + U(x,t)U_{x}(x,t) + \frac{1}{2}U_{xxx}(x,t) = 0,$$

$$0 < \alpha \le 1,$$
(12)

with initial condition

$$U(x,0) = 6\gamma^2 \operatorname{sech}^2(\gamma x). \tag{13}$$

Note that, this KdV equation have an exact solution when $\alpha = 1$ [52] as follows

$$U(x,t) = 6\gamma^2 \operatorname{sech}^2(\gamma x - 2\gamma^3 t). \tag{14}$$

Applying the MGMLFM to Eq.(12), and using Eq.(8), we assume

$$U(x,t) = F(x)E_{\alpha}(At^{\alpha}) = \sum_{n=0}^{\infty} F(x)A^{n} \frac{t^{n\alpha}}{\Gamma(n\alpha+1)},$$
(15)

where A is undetermined coefficient. From the initial condition (13), we have $F(x) = 6\gamma^2 \operatorname{sech}^2(\gamma x)$. By using Eq.(10), we obtain the linear part of Eq.(12) as the following



$$\mathcal{L}(U(x,t)) = \frac{1}{2} U_{xxx}(x,t) = \frac{1}{2} \frac{\partial^3 (6\gamma^2 \operatorname{sech}^2(\gamma x))}{\partial x^3}$$
$$= 4\gamma^3 U(x,0)(2 \tanh(\gamma x) \operatorname{sech}^2(\gamma x)$$
$$- \tanh^3(\gamma x)).$$

Similarly, the nonlinear part of Eq.(12) satisfies

$$\mathcal{N}(U(x,t)) = U(x,t)U_x(x,t)$$

$$= (6\gamma^2 \operatorname{sech}^2(\gamma x)) \frac{\partial (6\gamma^2 \operatorname{sech}^2(\gamma x))}{\partial x}$$

$$= U(x,0)(-12\gamma^3 \tanh(\gamma x) \operatorname{sech}^4(\gamma x)).$$

By using Eq.(9) we get

$$U(x,0) \sum_{n=0}^{\infty} (A^{n+1} - 12\gamma^3 \tanh(\gamma x) \operatorname{sech}^2(\gamma x) C^n \Gamma(n\alpha + 1)$$

$$+ 4\gamma^3 (2 \tanh(\gamma x) \operatorname{sech}^2(\gamma x)$$

$$- \tanh^3(\gamma x) A^n) \frac{t^{n\alpha}}{\Gamma(n\alpha + 1)} = 0, \quad (16)$$

where

$$C^{n} = \sum_{k=0}^{n} \frac{A^{k} A^{n-k}}{\Gamma(k\alpha+1)\Gamma((n-k)\alpha+1)}$$

Therefore, the recurrence relation will be as follows

$$A^{n+1} = 12\gamma^3 \tanh(\gamma x) \operatorname{sech}^2(\gamma x) C^n \Gamma(n\alpha + 1)$$
$$-4\gamma^3 (2 \tanh(\gamma x) \operatorname{sech}^2(\gamma x)$$
$$-\tanh^3(\gamma x)) A^n. \tag{17}$$

By substitute different values of n and doing some important calculation, we have

$$A^{0} = 1,$$

$$A^{1} = 4\gamma^{3} \tanh(\gamma x),$$

$$A^{2} = 16\gamma^{6} \tanh^{2}(\gamma x)(3 \operatorname{sech}^{2}(\gamma x) + 1),$$

$$A^{3} = 64\gamma^{9} \tanh^{3}(\gamma x)(9 \operatorname{sech}^{4}(\gamma x) + 6 \operatorname{sech}^{2}(\gamma x) + 3 \operatorname{sech}^{2}(\gamma x) \frac{\Gamma(2\alpha + 1)}{\Gamma(\alpha +)^{2}} + 1).$$

From Eq.(15), we obtain approximate solution in a series form as the following

$$U(x,t) = 6\gamma^{2} \operatorname{sech}^{2}(\gamma x) (1 + 4\gamma^{3} \tanh(\gamma x) \frac{t^{\alpha}}{\Gamma(\alpha + 1)} + 16\gamma^{6} \tanh^{2}(\gamma x) (3 \operatorname{sech}^{2}(\gamma x) + 1) \frac{t^{2\alpha}}{\Gamma(2\alpha + 1)} + 64\gamma^{9} \tanh^{3}(\gamma x) (9 \operatorname{sech}^{4}(\gamma x) + 6 \operatorname{sech}^{2}(\gamma x) + 3 \operatorname{sech}^{2}(\gamma x) \frac{\Gamma(2\alpha + 1)}{\Gamma(\alpha + 1)^{2}} + 1) \frac{t^{3\alpha}}{\Gamma(3\alpha + 1)} + \cdots).$$
(18)

In fact, these results satisfy the exact solution shown earlier in Eq.(14) when $\alpha = 1$, which means that the approximate solutions obtained by MGMLFM are rapidly convergent to the exact solutions and this is illustrated by the simulation presented in Section 5.

4.2 Time-Fractional KdVB Equation

We consider the time-fractional KdVB equation as follows

$${}^{C}_{0}D^{\alpha}_{t}U(x,t) + U(x,t)U_{x}(x,t) - \delta U_{xx}(x,t) + \frac{1}{2}U_{xxx}(x,t) = 0, \quad 0 < \alpha \le 1, (19)$$

with initial condition

$$U(x,0) = a\delta^2(2 - 2\tanh(\phi) + \mathrm{sech}^2(\phi)),$$
 (20)

where

$$a = \frac{6}{25}$$
 and $\phi = \frac{\delta x}{5}$.

This KdVB equation is exactly solvable when $\alpha = 1$ [52] and introduced as follows

$$U(x,t) = 2a\delta^2 - 2a\delta^2 \tanh(\omega) + 2a\delta^2 \operatorname{sech}^2(\omega), (21)$$

where $\omega=(\frac{-\delta x}{5}-\frac{2a\delta^3t}{5})$. Applying the MGMLFM to Eq.(19), and using Eq.(8), we assume

$$U(x,t) = \mathcal{W}(x)E_{\alpha}(At^{\alpha}) = \sum_{n=0}^{\infty} \mathcal{W}(x)A^{n} \frac{t^{n\alpha}}{\Gamma(n\alpha+1)}.$$
(22)

From Eq.(20), we have $\mathcal{W}(x) = a\delta^2(2-2\tanh(\phi)+\mathrm{sech}^2(\phi))$. By using Eq.(10), we obtain the linear part of Eq.(19) as the

$$\mathcal{L}(U(x,t)) = \frac{1}{2} U_{xxx}(x,t) - \delta U_{xx}(x,t)$$

$$= \frac{1}{2} \frac{\partial^3 (a\delta^2 (2 - 2\tanh(\phi) + \operatorname{sech}^2(\phi)))}{\partial x^3}$$

$$-\delta \frac{\partial^2 (a\delta^2 (2 - 2\tanh(\phi) + \operatorname{sech}^2(\phi)))}{\partial x^2}$$

$$= \frac{12a\delta^5}{125} \operatorname{sech}^2(\phi) (3\operatorname{sech}^2(\phi)$$

$$-2\tanh(\phi) - 2 + \tanh(\phi)\operatorname{sech}^2(\phi)).$$

Similarly, the nonlinear part of Eq.(19) satisfies

$$\mathcal{N}(U(x,t)) = U(x,t)U_x(x,t)$$

$$= (a\delta^2(2 - 2\tanh(\phi) + \mathrm{sech}^2(\phi)))$$

$$\frac{\partial(a\delta^2(2 - 2\tanh(\phi) + \mathrm{sech}^2(\phi)))}{\partial x}$$

$$= \frac{12a\delta^5}{125}\operatorname{sech}^2(\phi)(-3\operatorname{sech}^2(\phi))$$

$$- \tanh(\phi)\operatorname{sech}^2(\phi)).$$

By using Eq.(9) we get



$$\begin{split} \sum_{n=0}^{\infty} (\,a\delta^2(2-2\tanh(\phi)+\mathrm{sech}^2(\phi))A^{n+1} + \frac{12a\delta^5}{125}\,\mathrm{sech}^2(\phi)[(\,-3\,\mathrm{sech}^2(\phi)-\tanh(\phi)\,\mathrm{sech}^2(\phi))\,C^n\Gamma(n\alpha+1) \\ + (\,3\,\mathrm{sech}^2(\phi)-2\tanh(\phi)-2+\tanh(\phi)\,\mathrm{sech}^2(\phi))\,A^n])\,\frac{t^{n\alpha}}{\Gamma(n\alpha+1)} = 0, \end{split}$$

where

$$C^{n} = \sum_{k=0}^{n} \frac{A^{k} A^{n-k}}{\Gamma(k\alpha+1)\Gamma((n-k)\alpha+1)}.$$

Therefore, the recurrence relation will be as follows

$$A^{n+1} = \frac{12\delta^3}{125}\operatorname{sech}^2(\phi)$$

$$\frac{(3\operatorname{sech}^2(\phi) + \tanh(\phi)\operatorname{sech}^2(\phi))C^n\Gamma(n\alpha + 1) - (3\operatorname{sech}^2(\phi) - 2\tanh(\phi) - 2 + \tanh(\phi)\operatorname{sech}^2(\phi))A^n}{2 - 2\tanh(\phi) + \operatorname{sech}^2(\phi)}$$

By substitute different values of n and doing some important calculation, we have:

$$\begin{split} A^0 &= 1, \\ A^1 &= \frac{24\delta^3}{125} \operatorname{sech}^2(\phi) \big(\frac{\tanh(\phi) + 1}{2 - 2\tanh(\phi) + \operatorname{sech}^2(\phi)} \big), \\ A^2 &= \frac{12\delta^3}{125} \operatorname{sech}^2(\phi) \frac{(3\operatorname{sech}^2(\phi) + \tanh(\phi)\operatorname{sech}^2(\phi)) \, 2A^0A^1 - (3\operatorname{sech}^2(\phi) - 2\tanh(\phi) - 2 + \tanh(\phi)\operatorname{sech}^2(\phi)) \, A^1}{2 - 2\tanh(\phi) + \operatorname{sech}^2(\phi)}, \\ A^3 &= \frac{12\delta^3}{125} \operatorname{sech}^2(\phi) \frac{(3\operatorname{sech}^2(\phi) + \tanh(\phi)\operatorname{sech}^2(\phi)) \, C^2 - (3\operatorname{sech}^2(\phi) - 2\tanh(\phi) - 2 + \tanh(\phi)\operatorname{sech}^2(\phi)) \, A^2}{2 - 2\tanh(\phi) + \operatorname{sech}^2(\phi)}, \\ &: \end{split}$$

where

$$C^{2} = (2A^{0}A^{2} + (\frac{A^{1}}{\Gamma(\alpha+1)})^{2}\Gamma(2\alpha+1))$$

From Eq.(22), we obtain approximate solution as the following

$$U(x,t) = a\delta^{2}(2 - 2\tanh(\phi) + \mathrm{sech}^{2}(\phi))(A^{0} + A^{1}\frac{t^{\alpha}}{\Gamma(\alpha+1)} + A^{2}\frac{t^{2\alpha}}{\Gamma(2\alpha+1)} + A^{3}\frac{t^{3\alpha}}{\Gamma(3\alpha+1)} + \cdots).$$
 (23)

This approximate solution obtained by MGMLFM coincides with the exact solution given by Eq.(21) when $\alpha=1$. This confirms that the approximate solution obtained by the proposed method is rapidly convergent to the exact solutions and this is explained from the simulations presented in the following section.

5 Results and discussion

In this section, we introduce the simulation of our results. In Figure 1, we present the MGMLFM solution (18) of time-fractional KdV equation (12) when $\gamma=0.1$, with different values of $\alpha=1,0.9,0.8$ and comparing with the

exact solution given by Eq.(14). Figure 2 shows the absolute errors for time-fractional KdV when $\alpha=1$ among different values of γ and it becomes clear that whenever decreases the value of γ decreases the absolute errors. In Figure 3, we show the MGMLFM solution (23) of time-fractional KdVB equation (19) when $\delta=0.1$, with different values of $\alpha=1,0.9,0.8$ and comparing with the exact solution given by Eq.(21). Figure 4 shows the absolute errors for the time-fractional KdVB when $\alpha=1$ with different values of δ and it becomes clear that whenever δ very small the absolute errors decrease, hence we obtain an accurate approximation. These graphs prove that the solution obtained by the proposed method approach the exact solution when $\alpha \to 1$.



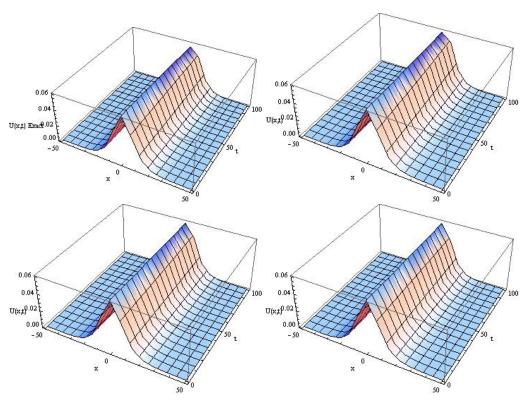


Fig. 1: The exact and approximate solutions obtained by MGMLFM of Eq.(12) for $\gamma=0.1$ with $\alpha=1,\,\alpha=0.9$ and $\alpha=0.8$, respectively.

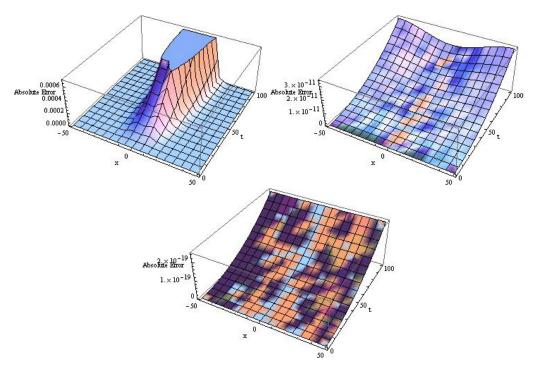


Fig. 2: The absolute error between the exact solution and MGMLF solution when $\alpha=1$ for KdV equation (12) when $\gamma=0.1,0.01$ and 0.001, respectively.

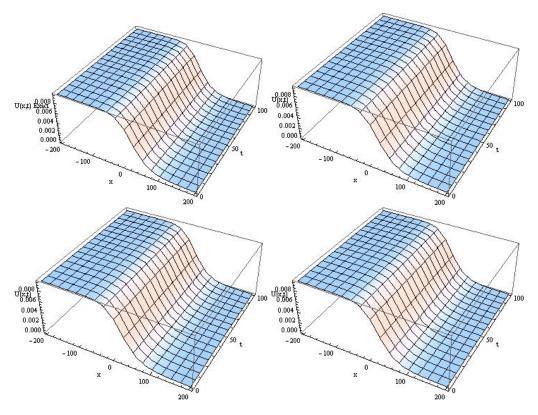


Fig. 3: The exact and approximate solutions obtained by MGMLFM of Eq.(19) for $\delta=0.1$ with $\alpha=1,\,\alpha=0.9$ and $\alpha=0.8$, respectively.

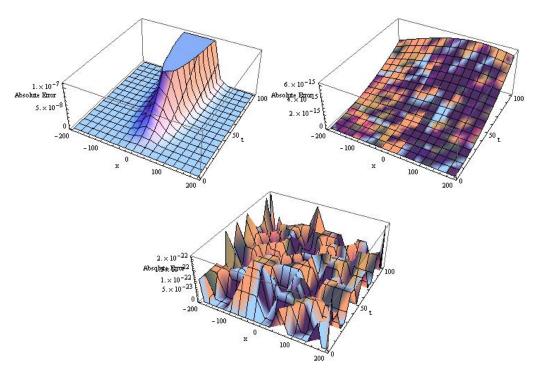


Fig. 4: The absolute error between the exact solution and MGMLF solution when $\alpha=1$ for KdVB equation (19) when $\delta=0.1,0.01$ and 0.001, respectively.



Table 1: Value of solutions obtained by MGMLFM, exact solution, and absolute errors of KdV equation (12) with different values of t and γ when $\alpha = 1$.

γ	x	t	MGMLFM solution	Exact solution	Absolute error
0.1	30	5	0.000603864	0.00060386	4.15823×10^{-9}
		15	0.000628414	0.000628375	3.91059×10^{-8}
		25	0.000653992	0.000653879	1.12544×10^{-7}
0.01	30	5	0.000549085	0.000549085	7.58348×10^{-14}
		15	0.000549092	0.000549092	6.82526×10^{-13}
		25	0.000549098	0.000549098	1.89594×10^{-12}
0.001	30	5	5.9946×10^{-6}	5.9946×10^{-6}	8.47033×10^{-22}
		15	5.9946×10^{-6}	5.9946×10^{-6}	4.23516×10^{-21}
		25	5.9946×10^{-6}	5.9946×10^{-6}	1.52466×10^{-20}

Table 2: Value of solutions obtained by MGMLFM, exact solution, and absolute errors of KdVB equation (19) with different values of t and δ when $\alpha = 1$.

δ	x	t	MGMLFM solution	Exact solution	Absolute error
0.1	30	5	0.00393247	0.00393247	1.1008×10^{-9}
		15	0.00393752	0.00393751	9.92569×10^{-9}
		25	0.00394259	0.00394256	2.76226×10^{-8}
0.01	30	5	0.0000690373	0.0000690373	1.52466×10^{-17}
		15	0.0000690373	0.0000690373	1.37219×10^{-16}
		25	0.0000690374	0.0000690374	3.81165×10^{-16}
0.001	30	5	7.17111×10^{-7}	7.17111×10^{-7}	1.05879×10^{-22}
		15	7.17111×10^{-7}	7.17111×10^{-7}	1.05879×10^{-22}
		25	7.17111×10^{-7}	7.17111×10^{-7}	0

6 Conclusion

In this article, the MGMLFM has been successfully employed to obtain approximate-analytical solutions of time-fractional KdV and KdVB equations, Eq.(12) and Eq.(19), respectively. The MGMLFM gives series solutions as in Eq.(18) and Eq.(23), which converge rapidly, and require less computational work, and provide highly accurate results when comparing approximate solutions when $\alpha = 1$ with recognized exact solutions which are shown in Eq.(14) and Eq.(21). Some simulation of the obtained results was presented in the forms of graphs and tables. As it turns out that the absolute error can be controlled through the parameters in the time-fractional KdV and KdVB equations. From the previous arguments, It becomes clear to us that placing into the practice of the MGMLFM method is very reliable, well-organized, easily and it relevant to solve further nonlinear FPDEs that have various applications in different areas of applied sciences. We used Mathematica software for computations and plotting the figures.

Conflict of Interest

The authors declare that there is no conflict of interest regarding the publication of this article.

References

- [1] D. J. Kordeweg, G. de Vries. On the change of form of long waves advancing in a rectangular channel, and a new type of long stationary wave. Philos. Mag. 39, 422-443, 1895.
- [2] F. u. Zuntao, L. Shida, L. Shikuo. New solutions to mkdv equation. Phys. Lett. A. 326(5), 364-374, 2004.
- [3] D. Kaya. An application for the higher order modified KdV equation by decomposition method. Commun. Nonlinear Sci. Numer. Simul. 10, 693-702, 2005.
- [4] A. M. Wazwaz. New sets of solitary wave solutions to the kdv, mkdv, and the generalized kdv equations. Commun. Nonlinear Sci. Numer. Simul. 13(2), 331-339, 2008.
- [5] S. R. Mousavian, H. Jafari, C. M. Khalique, S. A. Karimi. New exact-analytical solutions for the mkdv equation. TJMCS. 2(3), 413-416, 2011.
- [6] C. H. Su, C. S. Gardner. Korteweg-de Vries equation and generalizations III. Derivation of the Korteweg-de Vries equation and Burgers equation. J. Math. Phys. 10(3), 536-539, 1969.
- [7] Z. Feng. On travelling wave solutions of the Burgers-Korteweg-de Vries equation. Nonlinearity. 20, 343-356,
- [8] A. K. Gupta, S. S. Ray. On the solution of time-fractional KdV-Burgers equation using Petrov-Galerkin method for propagation of long wave in shallow water. Chaos. Soliton. Fract. 116, 376-380, 2018.
- [9] L. van Wijngaarden. On the motion of gas bubbles in a perfect fluid. Ann. Rev. Fluid Mech. 4, 369-73, 1972.
- [10] P. N. Hu. Collisional theory of shock and nonlinear waves in a plasma Phys. Fluids. 15, 854-64, 1972.
- [11] M. S. Ruderman. Method of derivation of the Kortewegde Vries-Burgers equation. J. Appl. Math. Mech. 39, 656-64,
- [12] T. Karahara. Weak nonlinear magneto-acoustic waves in a cold plasma in the presence of effective electron-ion collisions. J. Phys. Soc. Japan. 27, 1321-9, 1970.
- [13] S. D. Liu, S. K. Liu. KdV-Burgers equation modelling of turbulence. Sci. Sin. (Ser. A). 35, 576-86, 1992.
- [14] F. Feudel, H. Steudel, Nonexistence of prolongation structure for the Korteweg-de Vries-Burgers equation. Phys. Lett. A. 107, 5-8, 1985.
- [15] D. Kaya. An application for the higher order modified KdV equation by decomposition method. Comm. Nonlinear Sci. Numer. Simul. 10, 693-702, 2005.
- [16] Q. Wang. Numerical solution for fractional KdV-Burgers equation by Adomian decomposition method. Appl. Math. Comput. 182, 1048-55, 2006.
- [17] I. Dağ, Y. Dereli. Numerical solutions of KdV equation using radial basis function, Appl. Math. Model., 32, 535-546,
- [18] S. Haq, S. U. Islam, M. Uddin. A mesh-free method for the numerical solution of the KdV-Burgers equation. Appl. Math. Model. 33, 3442-9, 2009.
- [19] A.M. Wazwaz. The extended tanh method for new solitons solutions for many forms of the fifth-order KdV equations. Appl. Math. Comput. 184, 1002-1014, 2007.
- [20] B. Sahu, R. Roychoudhury. Travelling wave solution of Korteweg-de Vries-Burger's equation. Czech J. Phys. 53, 517-27, 2003.



- [21] S. Abbasbandy, F. S. Zakaria. Soliton solutions for the fifthorder KdV equation with the homotopy analysis method. Nonlinear Dyn., 51, 83–87, 2008.
- [22] Q. Wang. Homotopy perturbation method for fractional KdV–Burgers equation. Chaos. Soliton. Fract. 35, 843–50, 2008.
- [23] C. Koroglu. Exact and nonstandard finite difference schemes for the generalized KdV–Burgers equation. Advan. Difference Equat. 134, 1-21, 2020.
- [24] W. Cao, Y. Xu, Z. Zheng. Finite difference/collocation method for a Generalized time-fractional KdV equation. Appl.Sci. 8(1),1–15, 2018.
- [25] S. B. G. Karakoc, K. K. Ali. New exact solutions and numerical approximations of the generalized KdV equation. Comput. Meth. Differen. Equat. 2020, 1-21, 2020.
- [26] B. Saka, I. Dag. Quartic B-spline Galerkin approach to the numerical solutions of the KdVB equation. Appl. Math. Comput. 215, 746–58, 2009.
- [27] R. Hedli, A. Kadem. Exact Traveling Wave Solutions to the Fifth-order KdV Equation Using the Exponential Expansion Method. IAENG International J. Appl. Math., 50(1), 1-6, 2020.
- [28] H. Ahmad, T. A. Khan, S. W. Yao. An efficient approach for the numerical solution of fifth-order KdV equations. Open Mathematics, 18(1), 738-748, 2020.
- [29] S. Sahoo, S. Saha Ray. A new method for exact solutions of variant types of time fractional Korteweg-de-Vries equations in shallow water waves. Math. Methods Appl. Sci. 40, 106–14, 2017.
- [30] A. El-Ajou, O.A. Arqub, S. Momani. Approximate analytical solution of the nonlinear fractional KdV–Burgers equation: a new iterative algorithm. J. Comput. Phys. 293, 81–95, 2015.
- [31] M. A. Helal, M. S. Mehanna. A comparison between two different methods for solving KdV–Burgers equation. Chaos, Soliton. Fract. 28, 320–6, 2006.
- [32] M. A. Herzallah, A. M. El-Sayed, D. Baleanu. On the fractional-order diffusion-wave process. Rom. J. Phys., 55(3-4), 274-284, 2010.
- [33] M. Yavuz and E. Bonyah. New approaches to the fractional dynamics of schistosomiasis disease model, Physica A. 525, 373-393, 2019.
- [34] M. B. Riaz, M. A. Imran, K. Shabbir. Analytic solutions of Oldroyd-B fluid with fractional derivatives in a circular duct that applies a constant couple, Alex. Eng. J. 55, 3267-3275, 2016.
- [35] H. M. Ali, F. L. Pereira, S. Gama. A new approach to the Pontryagin maximum principle for nonlinear fractional optimal control problems, Math. Meth. Appl. Sci. 39(13), 3640-364, 2016.
- [36] R. Hilfer. Applications of Fractional Calculus in Physics, World Scientific, Singapore, 2000.
- [37] C. Park, M. M. Khater, A. H. Abdel-Aty, R. A. Attia, H. Rezazadeh, A. M. Zidan, A. B. Mohamed. Dynamical analysis of the nonlinear complex fractional emerging telecommunication model with higher–order dispersive cubic–quintic. Alex. Eng. J. 59, 1425-1433, 2020.
- [38] M. A. Abdou, S. Owyed, A. Abdel-Aty, B. M. Raffah, S. Abdel-Khalek. Optical soliton solutions for a space-time fractional perturbed nonlinear Schrödinger equation arising in quantum physics. Resul. Phys. 16, 102895, 2020.

- [39] S. Owyed, M. A. Abdou, A. H. Abdel-Aty, W. Alharbi, R. Nekhili. Numerical and approximate solutions for coupled time fractional nonlinear evolutions equations via reduced differential transform method. Chaos Soliton. Fract.131, 109474, 2020.
- [40] S. Owyed, M. A. Abdou, A. H. Abdel-Aty, A. A. Ibraheem, R. Nekhili, D. Baleanu. New optical soliton solutions of space-time fractional nonlinear dynamics of microtubules via three integration schemes. J. Intel. Fuzzy Sys., 38, 2859 – 2866, 2020.
- [41] S. Owyed, M. A. Abdou, A. Abdel-Aty, H. Dutta. Optical solitons solutions for perturbed time fractional nonlinear Schrodinger equation via two strategic algorithms. AIMS MATHEMATICS, 5(3), 2057-2070, 2020.
- [42] M. M.A. Khater, R. A. M. Attia, A. H. Abdel-Aty. Unidirectional propagation of long wave in dispersive media through new computational scheme. Math. Scie. Letters. 9, 15-21, 2020.
- [43] K. A. Gepreel, M. S. Mohamed. An optimal homotopy analysis method nonlinear fractional differential equation. JARDCS. 6(1), 1-10, 2014.
- [44] A. A. M. Arafa, S.Z. Rida, H. Mohamed. An application of the homotopy analysis method to the transient behavior of a biochemical reaction model. Inform. Sci. Lett. 31, 29–33, 2014.
- [45] M. S. Mohamed. Application of optimal HAM for solving the fractional order logistic equation. Appl. Comput. Math. **3**(1), 27-31, 2014.
- [46] A. A. M. Arafa, S. Z. Rida, H. Mohamed. Approximate analytical solutions of Schnakenberg systems by homotopy analysis method. Appl. Math. Mod. 36(10), 4789–4796, 2012.
- [47] A. A. M. Arafa, S. Z. Rida, H. Mohamed. Homotopy analysis method for solving biological population model. Commun. Theor. Phys. 56(5), 797–800, 2011.
- [48] S. Javeed, D. Baleanu, A. Waheed, M. S. Khan, H. Affan, Analysis of homotopy perturbation method for solving fractional order differential equations, Mathematics, 7(1), 1-14, 2019.
- [49] A. M. A. El-Sayed, A. Elsaid, I. L. El-Kalla, et al. A homotopy perturbation technique for solving partial differential equations of fractional order in finite domains, Appl. Math. Comput. 218, 8329-8340, 2012.
- [50] P. K. Gupta, M. Singh, Homotopy perturbation method for fractional fornberg-whitham equation. Comput. Math. Appl. 61, 250-254, 2011.
- [51] J. H. He. Application of homotopy perturbation method to nonlinear wave equation, Chaos. Soliton. Fract. 26, 695-700, 2005.
- [52] K. M. Saa, E. H. AL-Shareef, A. K. Alomari, D. Baleanu, J. F. Gómez-Aguilar. On exact solutions for time-fractional Korteweg-de Vries and Korteweg-de Vries-Burger's equations using homotopy analysis transform method. Chinese J. Phys. 63, 149-162, 2020.
- [53] M. Khan, M. A. Gondal, I. Hussain, S. K. Vanani. A new comparative study between homotopy analysis transform method and homotopy perturbation transform method on a semi infinite domain. Math. Compute. Model., 55(3-4), 1143-1150, 2012.
- [54] Y. Liu, H. Sun, X. Yin, B. Xin. A new Mittag-Leffler function undetermined coefficient method and its



- applications to fractional homogeneous partial differential equations, J. Nonlin. Sci. Appl. **10** (8), 4515-4523, 2017.
- [55] H. M. Ali, I. Ameen. Mittag-Leffler approximation for the solution of nonlinear systems of fractional partial differential equations. SYLWAN. 163(9), 17-32, 2019.
- [56] H. M. Ali, I. Ameen. An Efficient Approach for Solving Fractional Dynamics of a Predator-Prey System. Mod. Appl. Sci. 13 (11), 116-126, 2019.
- [57] H. M. Ali. New approximate solutions to fractional smoking model using the generalized Mittag-Leffler function method. Progr. Fract. Differ. Appl. 5(4), 319-326, 2019.
- [58] A.A.M. Arafa, S.Z. Rida, A.A. Mohammadein, H.M. Ali. solving nonlinear fractional differential equation by generalized Mittag-Leffler function method, Commun. Theor. Phys. 59(6), 661–663, 2013.
- [59] I. Podlubny. Fractional Differential Equations, Mathematics in Sciences and Engineering, Academic Press, San Diego, 1999.
- [60] D. Baleanu, K. Diethelm, E. Scalas and J. J. Trujillo. Fractional Calculus: Models and Numerical Methods, Series on Complexity, Nonlinearity and Chaos, Vol. 3, World Scientific Publishing Company, Singapore, New Jersey, London and Hong Kong, 2012.
- [61] A. Ghorbani. Beyond Adomian polynomials: He polynomials. Chaos. Soliton. Fract. 39, 1486-1492, 2009.
- [62] S. T. Mohyud-Din, M. A. Noor, K. I. Noor. Traveling wave solutions of seventh-order generalized KdV equations using he's polynomials. Int J Nonlin Sci. Num. 10, 227-233, 2009.
- [63] Y. Q. Liu. Approximate solutions of fractional nonlinear equations using homotopy perturbation transformation method. Abstr. Appl. Anal. 2012, Article ID 752869, 14 pages, 2012.