1227

# Spacelike Curves in the Lightlike Cone 

Nemat Abazari ${ }^{1, *}$, Martin Bohner ${ }^{2}$, Ilgın Sağer ${ }^{3}$ and Alireza Sedaghatdoost ${ }^{1}$<br>${ }^{1}$ Department of Mathematics and Applications, University of Mohaghegh Ardabili, Ardabil 56199-11367, Iran<br>${ }^{2}$ Department of Mathematics and Statistics, Missouri University of Science and Technology, Rolla, Missouri 65409-0020, USA<br>${ }^{3}$ Department of Mathematics and Computer Science, University of Missouri-St. Louis, St. Louis, Missouri 63121, USA

Received: 18 Mar. 2018, Revised: 29 Oct. 2018, Accepted: 29 Oct. 2018
Published online: 1 Nov. 2018


#### Abstract

In this paper, we study curves in the lightlike cone. First, we show that any curves in the lightlike cone are spacelike or lightlike, and then we characterize some curves with special cone curvature function in the 4,5 , and 6 -dimensional lightlike cone. Finally, we consider the relationship between Frenet curvature functions and cone curvature functions for a spacelike curve on the lightlike cone.


Keywords: Asymptotic frame, cone curvature, lightlike cone, spacelike curve

## 1 Introduction

In Euclidean space, we can consider the behavior of a curve by Frenet orthonormal frame and its Frenet curvatures. For example, if all of the Frenet curvatures of a curve in $\mathbb{E}^{n}$ are constant, then this curve is

$$
\begin{aligned}
\alpha(s)= & \left(a_{1} \sin \left(\alpha_{1} s\right), a_{1} \cos \left(\alpha_{1} s\right), \ldots\right. \\
& \left.\ldots, \alpha_{m} \sin \left(\alpha_{m} s\right), \alpha_{m} \cos \left(\alpha_{m} s\right), b s\right)
\end{aligned}
$$

for $n=2 m+1$ and

$$
\begin{aligned}
\alpha(s)= & \left(a_{1} \sin \left(\alpha_{1} s\right), a_{1} \cos \left(\alpha_{1} s\right), \ldots\right. \\
& \left.\ldots, \alpha_{m} \sin \left(\alpha_{m} s\right), \alpha_{m} \cos \left(\alpha_{m} s\right)\right)
\end{aligned}
$$

for $n=2 m$, see [3]. Also, S. Yılmaz and M. Turgut in [10] have defined vector products in Minkowski space-time, and by this way they calculate Frenet frames of all spacelike curves.

In the Lorentzian manifold, there are three type of curves, namely spacelike, timelike, and lightlike curves, and their Frenet equations are different, see [1,7].

Besides the Frenet orthonormal frame along a curve on a lightlike cone, an asymptotic orthonormal frame is very useful. Asymptotic orthonormal frames are applied in order to consider curves, surfaces, and hypersurfaces in the lightlike cones.
H. Liu in $[4,5]$ has considered curves in the lightlike cone $\mathbb{Q}^{n+1}$. For this consideration, he defined the
asymptotic orthonormal frame along a curve and cone curvature functions for such curve in $\mathbb{Q}^{n+1}$. Then he obtained some conformal invariants, and he classified curves with constant cone curvatures in $\mathbb{Q}^{2}$ and $\mathbb{Q}^{3}$. Also, in $\mathbb{Q}^{2}$, Liu established the relation between Frenet curvature and cone curvature, and he characterized the cone curvature function for a helix. We also remark that M. Bektaş and M. Külahci in [2] have obtained a characterization of spacelike curves in the 3-dimensional lightlike cone in terms of some differential equations.

In this article, we develop and generalize the results by H. Liu [4, 5]. The setup of this paper is as follows. After giving some preliminaries in Section 2, we show in Section 3 that any nonstraight line curve in $\mathbb{Q}^{n+1}$ is a spacelike curve. In Section 4, we characterize curves with constant cone curvature functions in $\mathbb{Q}^{4}, \mathbb{Q}^{5}$, and $\mathbb{Q}^{6}$. In Section 5, we give some relation between Frenet curvatures and cone curvature functions for a curve in $\mathbb{Q}^{3}$, and also we obtain cone curvature functions for a curve in $\mathbb{Q}^{3}$ such that the vectors $\alpha_{1}$ and $\alpha_{2}$ have constant angle with a constant vector $b$.

## 2 Preliminaries

Let $\mathbb{E}^{n}$ be $n$-dimensional Euclidean space. For two vectors $v=\left(v^{1}, \ldots, v^{n}\right), w=\left(w^{1}, \ldots, w^{n}\right)$ and an integer $q \in[0, n]$,

[^0]we define the bilinear form
$$
\langle v, w\rangle:=\sum_{i=1}^{n-q} v^{i} w^{i}-\sum_{i=n-q+1}^{n} v^{i} w^{i} .
$$

The resulting semi-Riemannian space is called Minkowski $n$-space, and $n=4$ is the simplest example of a relativistic space-time, see [7].

Definition 1. A vector $v \neq 0$ in $\mathbb{E}_{1}^{n}$ is called spacelike, timelike, or lightlike if $\langle v, v\rangle>0,\langle v, v\rangle<0$, or $\langle v, v\rangle=0$, respectively, and $v=0$ is spacelike.

Definition 2. The set of all lightlike vectors in $\mathbb{E}_{1}^{n}$ is called the lightlike cone and denoted by $\mathbb{Q}^{n-1}$.

In the lightlike cone $\mathbb{Q}^{n+1} \subset \mathbb{E}_{1}^{n+2}$, there are two orthonormal frame fields. One of them is a pseudo-Frenet orthonormal frame field and the other is an asymptotic orthonormal frame field, see [4].

Definition 3. A frame field $\left\{e_{1}, \ldots, e_{n}, e_{n+1}, e_{n+2}\right\}$ on $\mathbb{E}_{1}^{n+2}$ is called an asymptotic orthonormal frame field provided

$$
\begin{gathered}
\left\langle e_{n+1}, e_{n+1}\right\rangle=\left\langle e_{n+2}, e_{n+2}\right\rangle=0,\left\langle e_{n+1}, e_{n+2}\right\rangle=1 \\
\left\langle e_{n+1}, e_{i}\right\rangle=\left\langle e_{n+2}, e_{i}\right\rangle=0,\left\langle e_{i}, e_{j}\right\rangle=\delta_{i j}, i, j=1, \ldots, n
\end{gathered}
$$

Definition 4. A frame field $\left\{e_{1}, \ldots, e_{n}, e_{n+1}, e_{n+2}\right\}$ on $\mathbb{E}_{1}^{n+2}$ is called a pseudo-Frenet orthonormal frame field provided

$$
\begin{gathered}
\left\langle e_{n+1}, e_{n+1}\right\rangle=-\left\langle e_{n+2}, e_{n+2}\right\rangle=1,\left\langle e_{n+1}, e_{n+2}\right\rangle=0 \\
\left\langle e_{n+1}, e_{i}\right\rangle=\left\langle e_{n+2}, e_{i}\right\rangle=0,\left\langle e_{i}, e_{j}\right\rangle=\delta_{i j}, i, j=1, \ldots, n
\end{gathered}
$$

Definition 5. A curve $x$ in $\mathbb{E}_{1}^{n+2}$ is called a Frenet curve provided for all $t \in I$, the vector fields

$$
x(t), \dot{x}(t), \ddot{x}(t), \ldots, x^{(n)}(t), x^{(n+1)}(t)
$$

are linearly independent and the vector fields

$$
x(t), \dot{x}(t), \ddot{x}(t), \ldots, x^{(n+1)}(t), x^{(n+2)}(t)
$$

are linearly dependent, where $x^{(n)}(t)=\frac{\mathrm{d}^{n} x(t)}{\mathrm{d} t^{n}}$.
Definition 6. A curve $x: I \rightarrow \mathbb{E}_{1}^{n+2}$ is called spacelike, timelike, or lightlike provided $\dot{x}(t)$ is spacelike, timelike, or lightlike, respectively for all $t \in I$.
Definition 7. A spacelike or timelike curve $x: I \rightarrow \mathbb{E}_{1}^{n+2}$ is said to be parameterized by arc length provided

$$
\langle\dot{x}(s), \dot{x}(s)\rangle=1 \text { or }\langle\dot{x}(s), \dot{x}(s)\rangle=-1
$$

respectively.
Remark. In this article, all of the spacelike or timelike curves are parameterized by arclength denoted by $s$, and $x^{\prime}(s):=\frac{\mathrm{d} x(s)}{\mathrm{d} s}$.

In [4], H. Liu defined an asymptotic orthonormal frame field for a given curve in the $\mathbb{Q}^{n+1}$ as follows. Let $x: I \rightarrow$ $\mathbb{Q}^{n+1} \subset \mathbb{E}_{1}^{n+2}$ be a curve. We choose the null vector field $y(s)$ and the spacelike normal space $V^{n-1}$ of the curve $x$ such that they satisfy

$$
\begin{gathered}
\langle x(s), y(s)\rangle=1 \\
\langle x(s), x(s)\rangle=\langle y(s), y(s)\rangle=\left\langle x^{\prime}(s), y(s)\right\rangle=0 \\
V^{n-1}=\left(\operatorname{span}_{\mathbb{R}}\left\{x, y, x^{\prime}\right\}\right)^{\perp} \\
\operatorname{span}_{\mathbb{R}}\left\{x, y, x^{\prime}, V^{n-1}\right\}=\mathbb{E}_{1}^{n+2}
\end{gathered}
$$

Therefore, by choosing suitable orthonormal vector fields $\alpha_{2}(s), \alpha_{3}(s), \ldots, \alpha_{n}(s) \in V^{n-1}$, we have the Frenet formulas

$$
\begin{align*}
x^{\prime}(s) & =\alpha_{1}(s) \\
\alpha_{1}^{\prime}(s) & =\kappa_{1}(s) x(s)-y(s)+\tau_{1}(s) \alpha_{2}(s) \\
\alpha_{2}^{\prime}(s)= & \kappa_{2}(s) x(s)-\tau_{1}(s) \alpha_{1}(s)+\tau_{2}(s) \alpha_{3}(s) \\
& \vdots \\
\alpha_{i}^{\prime}(s)= & \kappa_{i}(s) x(s)-\tau_{i-1}(s) \alpha_{i-1}(s)+\tau_{i}(s) \alpha_{i+1}(s)  \tag{1}\\
& \vdots \\
\alpha_{n}^{\prime}(s)= & \kappa_{n}(s) x(s)-\tau_{n-1}(s) \alpha_{n-1}(s) \\
y^{\prime}(s)= & -\sum_{i=1}^{n} \kappa_{i}(s) \alpha_{i}(s)
\end{align*}
$$

where $\quad\left\{x(s), y(s), x^{\prime}(s), \alpha_{2}(s), \alpha_{3}(s), \ldots, \alpha_{n}(s)\right\} \quad$ is $\quad$ an asymptotic orthonormal frame field, called the asymptotic orthonormal frame on $\mathbb{E}_{1}^{n+2}$ along the curve $x$ in $\mathbb{Q}^{n+1}$. The functions $\kappa_{i}=\left\langle\alpha_{i}^{\prime}, y\right\rangle, \quad i=1, \ldots, n$ and $\tau_{i}=\left\langle\alpha_{i}^{\prime}, \alpha_{i+1}\right\rangle, i=1, \ldots, n-1$, are called cone curvature functions of the curve $x$.
Proposition 1(see [5]). Let $x: I \rightarrow \mathbb{Q}^{n+1} \subset \mathbb{E}_{1}^{n+2}$ be a spacelike curve and put

$$
\begin{equation*}
y(s):=-x^{\prime \prime}(s)-\frac{1}{2}\left\langle x^{\prime \prime}(s), x^{\prime \prime}(s)\right\rangle x(s) \tag{2}
\end{equation*}
$$

Then $\tau_{i}=0, i=1, \ldots, n$.

## 3 Nonstraight Line Curves in the Lightlike Cone $\mathbb{Q}^{n+1}$

In Euclidean space, a regular curve is a curve which has nonzero velocity vector. In the Minkowski space $\mathbb{E}_{1}^{3}$, any timelike (lightlike) curve is regular. Also, if a curve $x: I \rightarrow$ $\mathbb{E}_{1}^{3}$ is regular in $s_{0}$, then, by continuity, $x$ is also regular in a neighborhood of $s_{0}$, see [6]. Similarly to this, we can prove the following.

Proposition 2. Any timelike (lightlike) curve $x: I \rightarrow \mathbb{E}_{1}^{n+2}$ (with arbitrary parameter) is regular.

Proof. Assume that the curve is timelike. We write

$$
x(t)=\left(x_{1}(t), \ldots, x_{n+1}(t), x_{n+2}(t)\right)
$$

where $x_{i}$ are differentiable functions on $I$. In this case, we have

$$
\langle\dot{x}(t), \dot{x}(t)\rangle=\dot{x}_{1}^{2}(t)+\ldots+\dot{x}_{n+1}^{2}(t)-\dot{x}_{n+2}^{2}(t)<0 .
$$

In particular, $\dot{x}_{n+2}(t) \neq 0$, i.e., $x$ is regular. On the other hand, if the curve is lightlike, we have $\dot{x}_{n+2}(t) \neq 0$ again since, on the contrary, $\dot{x}_{i}(t)=0$ and $\dot{x}(t)=0$. But this means that the curve is spacelike.

Lemma 1. Let $x: I \rightarrow \mathbb{Q}^{n+1} \subset \mathbb{E}_{1}^{n+2}$ be a curve. Then $x$ is lightlike if and only if $x$ is a straight line.
Proof. Let $\langle x, x\rangle=0$ and $\langle\dot{x}, \dot{x}\rangle=0$, so

$$
\begin{align*}
& x_{n+2}^{2}=x_{1}^{2}+\ldots+x_{n+1}^{2},  \tag{3}\\
& \dot{x}_{n+2}^{2}=\dot{x}_{1}^{2}+\ldots+\dot{x}_{n+1}^{2}
\end{align*}
$$

Differentiation of the first equation in (3) yields that

$$
\begin{align*}
x_{n+2} \dot{x}_{n+2} & =x_{1} \dot{x}_{1}+\ldots+x_{n+1} \dot{x}_{n+1} \\
\left(x_{n+2}\right)^{2}\left(\dot{x}_{n+2}\right)^{2} & =\left(x_{1} \dot{x}_{1}+\ldots+x_{n+1} \dot{x}_{n+1}\right)^{2} \tag{4}
\end{align*}
$$

By substituting (3) into (4) and after some calculations, we conclude that

$$
\sum_{i, j=1}^{n+1}\left(x_{i} \dot{x}_{j}-x_{j} \dot{x}_{i}\right)^{2}=0
$$

Thus,

$$
\frac{\dot{x}_{i}}{x_{i}}=\frac{\dot{x}_{1}}{x_{1}}, i=1, \ldots, n+1
$$

Finally, we have

$$
\begin{aligned}
x_{i}(s) & =A_{i} x_{1}(s), i=1, \ldots, n+1 \\
x_{n+2}(s) & = \pm x_{1}(s) \sqrt{1+A_{1}^{2}+\ldots+A_{n+1}^{2}}
\end{aligned}
$$

where $A_{i}$ is some real constant. Thus, $x(s)=\vec{A} x_{1}(s)$ is a straight line with the real differentiable function $x_{1}$ and constant lightlike velocity vector $\vec{A}$. On the other hand, the converse statement is trivial.

Lemma 2. Let $x: I \rightarrow \mathbb{E}_{1}^{n+2}$ be a timelike curve. Then $x$ is not lying in $\mathbb{Q}^{n+1}$.

Proof. Assume that $x$ is in $\mathbb{Q}^{n+1}$. Then

$$
\begin{equation*}
x_{n+2}^{2}=x_{1}^{2}+\ldots+x_{n+1}^{2} \tag{5}
\end{equation*}
$$

so that

$$
\begin{equation*}
\sum_{i=1}^{n+1} x_{i} x_{i}^{\prime}=x_{n+2} x_{n+2}^{\prime} \tag{6}
\end{equation*}
$$

Since $x$ is timelike, we get

$$
\begin{align*}
&\left(x_{1}^{2}+\ldots+x_{n+1}^{2}\right)\left(x_{1}^{\prime 2}+\ldots+x_{n+1}^{\prime}{ }^{2}\right)-x_{n+2}^{2} x_{n+2}^{\prime} \\
&=-x_{n+2}^{2} \tag{7}
\end{align*}
$$

If we replace (6) in (7), then we obtain

$$
\sum_{i, j=1}^{n+1}\left(x_{i} x_{j}^{\prime}-x_{j} x_{i}^{\prime}\right)^{2}=-x_{n+2}^{2}
$$

Hence $x_{n+2}=0$, and (5) yields

$$
x_{i}(s)=0, i=1, \ldots, n+1
$$

So $x(s)=0$, which is a contradiction.
Lemmas 1 and 2 yield the following theorem.
Theorem 1. If $x: I \rightarrow \mathbb{Q}^{n+1} \subset \mathbb{E}_{1}^{n+2}$ is a regular curve, then $x$ is a nonstraight line if and only if $x$ is a spacelike curve.
Proof. Let $x$ be a nonstraight line curve in $\mathbb{Q}^{n+1}$. Then, by Lemma 1, this curve is not a lightlike curve, and, by Lemma 2, this curve is not a timelike curve. Therefore, it is a spacelike curve. Conversely, if the curve is spacelike and a straight line, then $x(s)=\vec{A} \tilde{x}(s)$ such that $\tilde{x}(s)$ is a real differentiable function and $\vec{A}$ is a lightlike vector as $x$ is lightlike, a contradiction.
Remark. For the rest of this article, we assume that the curve $x$ is a spacelike curve parameterized by arc length.

## 4 Curves in the Lightlike Cones $\mathbb{Q}^{4}, \mathbb{Q}^{5}$, and $\mathbb{Q}^{6}$

H. Liu in [4, Theorems 2.3 and 3.1] has classified all curves with constant cone curvature functions on $\mathbb{Q}^{2}$ and $\mathbb{Q}^{3}$. These curves are solutions of special differential equations. Similarly to these two theorems, we obtain constant cone curvature curves in the lightlike cones $\mathbb{Q}^{4}$, $\mathbb{Q}^{5}$, and $\mathbb{Q}^{6}$.
Theorem 2. Let $x: I \rightarrow \mathbb{Q}^{4} \subset \mathbb{E}_{1}^{5}$ be a curve in the lightlike cone $\mathbb{Q}^{4}$. If the cone curvature functions of the curve are constant, then the curve satisfies the differential equation

$$
\begin{align*}
x^{(5)}+\left(\tau_{1}^{2}+\right. & \left.\tau_{2}^{2}-2 \kappa_{1}\right) x^{\prime \prime \prime} \\
& -\left(\kappa_{2}^{2}+\kappa_{3}^{2}+2 \kappa_{3} \tau_{1} \tau_{2}+2 \kappa_{1} \tau_{2}^{2}\right) x^{\prime}=0 \tag{8}
\end{align*}
$$

Proof. The Frenet formulas (1) for this curve are

$$
\begin{align*}
x^{\prime}(s) & =\alpha_{1}(s) \\
\alpha_{1}^{\prime}(s) & =\kappa_{1}(s) x(s)-y(s)+\tau_{1}(s) \alpha_{2}(s) \\
\alpha_{2}^{\prime}(s) & =\kappa_{2}(s) x(s)-\tau_{1}(s) \alpha_{1}(s)+\tau_{2}(s) \alpha_{3}(s)  \tag{9}\\
\alpha_{3}^{\prime}(s) & =\kappa_{3}(s) x(s)-\tau_{2}(s) \alpha_{2}(s) \\
y^{\prime}(s) & =-\kappa_{1}(s) \alpha_{1}(s)-\kappa_{2}(s) \alpha_{2}(s)-\kappa_{3}(s) \alpha_{3}(s) .
\end{align*}
$$

From (9), we obtain the other derivatives of $x$ as

$$
\begin{align*}
x^{\prime \prime \prime}= & \tau_{1} \kappa_{2} x+\left(2 \kappa_{1}-\tau_{1}^{2}\right) x^{\prime}+\kappa_{2} \alpha_{2} \\
& +\left(\kappa_{3}+\tau_{1} \tau_{2}\right) \alpha_{3}  \tag{10}\\
x^{(4)}= & \left(\kappa_{2}^{2}+\kappa_{3}^{2}+\kappa_{3} \tau_{1} \tau_{2}\right) x \\
& +\left(2 \kappa_{1}-\tau_{1}^{2}\right) x^{\prime \prime}-\left(\kappa_{3} \tau_{2}+\tau_{1} \tau_{2}^{2}\right) \alpha_{2}+\kappa_{2} \tau_{2} \alpha_{3} \\
x^{(5)}= & -\kappa_{2} \tau_{1} \tau_{2}^{2} x+\left(\kappa_{2}^{2}+\kappa_{3}^{2}+2 \kappa_{3} \tau_{1} \tau_{2}+\tau_{1}^{2} \tau_{2}^{2}\right) x^{\prime} \\
& +\left(2 \kappa_{1}-\tau_{1}^{2}\right) x^{\prime \prime \prime}-\kappa_{2} \tau_{2}^{2} \alpha_{2} \\
& -\tau_{2}^{2}\left(\kappa_{3}+\tau_{1} \tau_{2}\right) \alpha_{3} \tag{11}
\end{align*}
$$

If we multiply (10) by $\tau_{2}^{2}$ and add the resulting equation to (11), then we obtain (8).

Corollary 1. If $x: I \rightarrow \mathbb{Q}^{4} \subset \mathbb{E}_{1}^{5}$ is a curve with constant cone curvature in the lightlike cone $\mathbb{Q}^{4}$, then the following cases hold.
i) Assume $\kappa_{3}=\tau_{2}=0$.
$i$-i) If $\kappa_{2}=0$, then

$$
x(s)=a_{1} s^{2}+a_{2} s+a_{3}
$$

provided $\lambda=0$,

$$
a_{1} \sinh (\sqrt{2 \lambda}) s+a_{2} \cosh (\sqrt{2 \lambda}) s+a_{3}
$$

provided $\lambda>0$, and

$$
a_{1} \sin (\sqrt{-2 \lambda}) s+a_{2} \cos (\sqrt{-2 \lambda}) s+a_{3}
$$

provided $\lambda<0$, where $\lambda:=2 \kappa_{1}-\tau_{1}^{2}$.
$i-i i)$ If $\kappa_{2} \neq 0$, then

$$
\begin{aligned}
x(s)= & a_{1} \sinh (\mu s)+a_{2} \cosh (\mu s)+a_{3} \sin (v s) \\
& +a_{4} \cos (v s)+a_{5}
\end{aligned}
$$

where $\pm \mu$ and $\pm \mathrm{iv}$ are the real and imaginary roots of the equation

$$
t^{4}-\left(2 \kappa_{1}-\tau_{1}^{2}\right) t^{2}-\kappa_{2}^{2}=0
$$

ii) If $\kappa_{3}=0$ and $\tau_{2} \neq 0$, then

$$
\begin{aligned}
x(s)= & a_{1} \sinh (\mu s)+a_{2} \cosh (\mu s)+a_{3} \sin (v s) \\
& +a_{4} \cos (v s)+a_{5},
\end{aligned}
$$

where $\pm \mu$ and $\pm \mathrm{i} v$ are the real and imaginary roots of the equation

$$
t^{4}+\left(\tau_{1}^{2}+\tau_{2}^{2}-2 \kappa_{1}\right) t^{2}-\left(\kappa_{2}^{2}+2 \kappa_{1} \tau_{2}^{2}\right)=0
$$

iii) Assume $\kappa_{3} \neq 0$. Consider the equation

$$
\begin{align*}
t^{4}+\left(\tau_{1}^{2}\right. & \left.+\tau_{2}^{2}-2 \kappa_{1}\right) t^{2} \\
& -\left(\kappa_{2}^{2}+\kappa_{3}^{2}+2 \kappa_{3} \tau_{1} \tau_{2}+2 \kappa_{1} \tau_{2}^{2}\right)=0 \tag{12}
\end{align*}
$$

iii-i) If the roots of (12) are $\pm \mu$ and $\pm \mathrm{i} v$, then

$$
\begin{align*}
x(s)= & a_{1} \sinh (\mu s)+a_{2} \cosh (\mu s)  \tag{13}\\
& +a_{3} \sin (v s)+a_{4} \cos (v s)+a_{5} .
\end{align*}
$$

iii-ii) If the roots of (12) are $\pm \mu$ and $\pm v$, then

$$
\begin{align*}
x(s)= & a_{1} \sin (\mu s)+a_{2} \cos (\mu s)+a_{3} \sin (v s) \\
& +a_{4} \cos (v s)+a_{5} . \tag{14}
\end{align*}
$$

iii-iii) If the roots of (12) are $\pm \mathrm{i} \mu$ and $\pm \mathrm{i} v$, then

$$
\begin{align*}
x(s)= & a_{1} \sinh (\mu s)+a_{2} \cosh (\mu s) \\
& +a_{3} \sinh (v s)+a_{4} \cosh (v s)+a_{5} . \tag{15}
\end{align*}
$$

iii-vi) If the roots of (12) are $\pm \mu \pm \mathrm{i} v$, then

$$
\begin{align*}
x(s)= & a_{1} \sinh (\mu s) \sin (v s) \\
& +a_{2} \cosh (\mu s) \sin (v s)  \tag{16}\\
& +a_{3} \sinh (\mu s) \cos (v s) \\
& +a_{4} \cosh (\mu s) \cos (v s)+a_{5},
\end{align*}
$$

Here, $a_{i} \in \mathbb{E}_{1}^{4}, i \in\{1,2,3,4,5\}$, are suitable constant vectors.

Proof. First, we prove i). If we set $\kappa_{3}=\tau_{2}=0$ in (9), then $\alpha_{3}$ is constant, and thus [4, Theorem 3.1] yields the statement. Next, we prove ii). The discriminant of the quadratic equation

$$
\begin{equation*}
t^{2}+\left(\tau_{1}^{2}+\tau_{2}^{2}-2 \kappa_{1}\right) t-\left(\kappa_{2}^{2}+2 \kappa_{1} \tau_{2}^{2}\right)=0 \tag{17}
\end{equation*}
$$

is
$\Delta= \begin{cases}\left(2 \kappa_{1}-\tau_{1}^{2}\right)^{2}+\tau_{2}^{4}+2 \tau_{1}^{2} \tau_{2}^{2}+4 \kappa_{1} \tau_{2}^{2}+4 \kappa_{2}^{2}, & \kappa_{1} \geq 0, \\ \left(2 \kappa_{1}-\tau_{2}^{2}\right)^{2}+\tau_{1}^{4}+2 \tau_{1}^{2} \tau_{2}^{2}-4 \kappa_{1} \tau_{1}^{2}+4 \kappa_{2}^{2}, & \kappa_{1}<0 .\end{cases}$
Since $\tau_{2} \neq 0$, we have $\Delta>0$, and thus (17) has two real roots with different signs. If we set $\tilde{x}:=x^{\prime}$ in (8), then we obtain a differential equation of order 4 with constant coefficients. After an integration, the statement is proved. Finally, we prove iii). If we set $\tilde{x}:=x^{\prime}$ in (8), then we obtain a differential equation of order 4 with constant coefficients. In general, the solution of such differential equations are $\exp (\lambda s)$, where $\lambda$ is a root of (12). If $\lambda$ is real, then the solution is

$$
a_{1} \sinh (\lambda s)+a_{2} \cosh (\lambda s)
$$

while if $\lambda$ is imaginary, then the solution is

$$
a_{3} \sin (i \lambda s)+a_{4} \cos (i \lambda s) .
$$

If $\lambda=\mu+\mathrm{i} v$ is a complex number, then

$$
\begin{aligned}
\tilde{x}(s)= & a_{1} \sinh (\mu s) \sin (v s)+a_{2} \cosh (\mu s) \sin (v s) \\
& +a_{3} \sinh (\mu s) \cos (v s)+a_{4} \cosh (\mu s) \cos (v s) .
\end{aligned}
$$

An integration yields that the curve is in the forms (13), (14), (15), or (16).

Theorem 3. Let $x: I \rightarrow \mathbb{Q}^{5} \subset \mathbb{E}_{1}^{6}$ be a curve in the lightlike cone $\mathbb{Q}^{5}$. If the cone curvature functions of the curve are constant, then the curve satisfies the differential equation

$$
\begin{align*}
x^{(6)}+\left(\tau_{1}^{2}+\tau_{2}^{2}+\tau_{3}^{2}-2 \kappa_{1}\right) x^{(4)} & \\
-\left(\kappa_{2}^{2}+\kappa_{4}^{2}+\left(2 \kappa_{1}-\tau_{1}^{2}\right)\left(\tau_{2}^{2}\right.\right. & \left.\left.+\tau_{3}^{2}\right)+\left(\kappa_{3}+\tau_{1} \tau_{2}\right)^{2}\right) x^{\prime \prime} \\
& =\left(\kappa_{4} \tau_{2}+\kappa_{2} \tau_{3}\right)^{2} x . \tag{18}
\end{align*}
$$

Proof. The Frenet formulas (1) for this curve are

$$
\begin{align*}
x^{\prime}(s)= & \alpha_{1}(s), \\
\alpha_{1}^{\prime}(s)= & \kappa_{1}(s) x(s)-y(s)+\tau_{1}(s) \alpha_{2}(s), \\
\alpha_{2}^{\prime}(s)= & \kappa_{2}(s) x(s)-\tau_{1}(s) \alpha_{1}(s)+\tau_{2}(s) \alpha_{3}(s), \\
\alpha_{3}^{\prime}(s) & =\kappa_{3}(s) x(s)-\tau_{2}(s) \alpha_{2}(s)+\tau_{3}(s) \alpha_{4}(s),  \tag{19}\\
\alpha_{4}^{\prime}(s) & =\kappa_{4}(s) x(s)-\tau_{3}(s) \alpha_{3}(s) \\
y^{\prime}(s) & =-\kappa_{1}(s) \alpha_{1}(s)-\kappa_{2}(s) \alpha_{2}(s) \\
& -\kappa_{3}(s) \alpha_{3}(s)-\kappa_{4}(s) \alpha_{4}(s) .
\end{align*}
$$

From (19), we obtain the other derivatives of $x$ as

$$
\begin{align*}
x^{\prime \prime \prime}= & \tau_{1} \kappa_{2} x+\left(2 \kappa_{1}-\tau_{1}^{2}\right) x^{\prime}+\kappa_{2} \alpha_{2} \\
& +\left(\kappa_{3}+\tau_{1} \tau_{2}\right) \alpha_{3}+\kappa_{4} \alpha_{4}  \tag{20}\\
x^{(4)}= & \left(\kappa_{2}^{2}+\kappa_{3}^{2}+\kappa_{4}^{2}+\kappa_{3} \tau_{1} \tau_{2}\right) x \\
& +\left(2 \kappa_{1}-\tau_{1}^{2}\right) x^{\prime \prime}-\left(\kappa_{3} \tau_{2}+\tau_{1} \tau_{2}^{2}\right) \alpha_{2} \\
& +\left(\kappa_{2} \tau_{2}-\kappa_{4} \tau_{3}\right) \alpha_{3}+\left(\kappa_{3} \tau_{3}+\tau_{1} \tau_{2} \tau_{3}\right) \alpha_{4}  \tag{21}\\
x^{(5)}= & \left(\kappa_{4} \tau_{1} \tau_{2} \tau_{3}-\kappa_{2} \tau_{1} \tau_{2}^{2}\right) x \\
& +\left(\kappa_{2}^{2}+\kappa_{3}^{2}+\kappa_{4}^{2}+\tau_{1}^{2} \tau_{2}^{2}+2 \kappa_{3} \tau_{1} \tau_{2}\right) x^{\prime} \\
& +\left(2 \kappa_{1}-\tau_{1}^{2}\right) x^{\prime \prime \prime}+\left(\kappa_{4} \tau_{2} \tau_{3}-\kappa_{2} \tau_{2}^{2}\right) \alpha_{2} \\
& -\left(\kappa_{3} \tau_{2}^{2}+\tau_{1} \tau_{2}^{3}+\kappa_{3} \tau_{3}^{2}+\tau_{1} \tau_{2} \tau_{3}^{2}\right) \alpha_{3} \\
& +\left(\kappa_{2} \tau_{2} \tau_{3}-\kappa_{4} \tau_{3}^{2}\right) \alpha_{4} \\
x^{(6)}= & \left(2 \kappa_{2} \kappa_{4} \tau_{2} \tau_{3}-\tau_{2}^{2}\left(\kappa_{2}^{2}+\kappa_{3}^{2}+\kappa_{3} \tau_{1} \tau_{2}\right)\right. \\
& \left.-\tau_{3}^{2}\left(\kappa_{3}^{2}+\kappa_{4}^{2}+\kappa_{3} \tau_{1} \tau_{2}\right)\right) x \\
& +\left(\kappa_{2}^{2}+\kappa_{3}^{2}+\kappa_{4}^{2}+\tau_{1}^{2} \tau_{2}^{2}+2 \kappa_{3} \tau_{1} \tau_{2}\right) x^{\prime \prime} \\
& +\left(2 \kappa_{1}-\tau_{1}^{2}\right) x^{(4)}+\left(\tau_{2}^{2}+\tau_{3}^{2}\right)\left(\kappa_{3} \tau_{2}+\tau_{1} \tau_{2}^{2}\right) \alpha_{2} \\
& +\left(\tau_{2}^{2}+\tau_{3}^{2}\right)\left(\kappa_{4} \tau_{3}-\kappa_{2} \tau_{2}\right) \alpha_{3} \\
& -\left(\tau_{2}^{2}+\tau_{3}^{2}\right)\left(\kappa_{3} \tau_{3}+\tau_{1} \tau_{2} \tau_{3}\right) \alpha_{4} \tag{22}
\end{align*}
$$

If we multiply (21) by $\left(\tau_{2}^{2}+\tau_{3}^{2}\right)$ and add the resulting equation to (22), then we obtain (18).
Corollary 2. Let $x: I \rightarrow \mathbb{Q}^{5} \subset \mathbb{E}_{1}^{6}$ be a curve with constant cone curvature in the lightlike cone $\mathbb{Q}^{5}$. Assume

$$
\kappa_{4} \tau_{2}+\kappa_{2} \tau_{3} \neq 0
$$

## Consider the equation

$$
\begin{align*}
& t^{3}+\left(\tau_{1}^{2}+\tau_{2}^{2}+\tau_{3}^{2}-2 \kappa_{1}\right) t^{2} \\
&-\left(\kappa_{2}^{2}+\kappa_{4}^{2}+\left(2 \kappa_{1}-\tau_{1}^{2}\right)\right.\left.\left(\tau_{2}^{2}+\tau_{3}^{2}\right)+\left(\kappa_{3}+\tau_{1} \tau_{2}\right)^{2}\right) t \\
&-\left(\kappa_{4} \tau_{2}+\kappa_{2} \tau_{3}\right)^{2}=0 \tag{23}
\end{align*}
$$

i) If the roots of (23) are $\lambda^{2}$ and $(\mu \pm \mathrm{i} v)^{2}$, then

$$
\begin{aligned}
x(s)= & a_{1} \sinh (\mu s) \sin (v s)+a_{2} \cosh (\mu s) \sin (v s) \\
& +a_{3} \sinh (\mu s) \cos (v s)+a_{4} \cosh (\mu s) \cos (v s) \\
& +a_{5} \sinh (\lambda s)+a_{6} \cosh (\lambda s) .
\end{aligned}
$$

ii) If the roots of (23) are $\lambda^{2}, \mu^{2}$, and $v^{2}$, then

$$
\begin{aligned}
x(s)= & a_{1} \sinh (\mu s)+a_{2} \cosh (\mu s)+a_{3} \sinh (v s) \\
& +a_{4} \cosh (v s)+a_{5} \sinh (\lambda s)+a_{6} \cosh (\lambda s) .
\end{aligned}
$$

iii) If the roots of (23) are $\lambda^{2}$, $(\mathrm{i} \mu)^{2}$, and (iv $)^{2}$, then

$$
\begin{aligned}
x(s)= & a_{1} \sin (\mu s)+a_{2} \cos (\mu s)+a_{3} \sin (v s) \\
& +a_{4} \cos (v s)+a_{5} \sinh (\lambda s)+a_{6} \cosh (\lambda s) .
\end{aligned}
$$

Here, $a_{i} \in \mathbb{E}_{1}^{6}, i \in\{1,2,3,4,5,6\}$, are suitable constant vectors.

Proof. Denote the left-hand side of (23) by $f(t)$. Then $f(0)<0$ and

$$
\begin{aligned}
f^{\prime}(t)= & 3 t^{2}+2\left(\tau_{1}^{2}+\tau_{2}^{2}+\tau_{3}^{2}-2 \kappa_{1}\right) t \\
& -\left(\kappa_{2}^{2}+\kappa_{4}^{2}+\left(2 \kappa_{1}-\tau_{1}^{2}\right)\left(\tau_{2}^{2}+\tau_{3}^{2}\right)\right. \\
& \left.+\left(\kappa_{3}+\tau_{1} \tau_{2}\right)^{2}\right), \\
\Delta= & 3\left(\tau_{2}^{2}+\tau_{3}^{2}\right)^{2}+\left(\tau_{2}^{2}+\tau_{3}^{2}+2\left(2 \kappa_{1}-\tau_{1}^{2}\right)\right)^{2} \\
& +12\left(\kappa_{2}^{2}+\kappa_{4}^{2}+\left(\kappa_{3}+\tau_{1} \tau_{2}\right)^{2}\right)>0 .
\end{aligned}
$$

Thus, (23) has at least one positive root. The other roots are one of the following. In i) both are $(\mu \pm \mathrm{i} v)^{2}$. In ii), both of them are positive, i.e., they are $\mu^{2}$ and $v^{2}$. In iii), both of them are negative, i.e., they $(\mathrm{i} \mu)^{2}$ and $(\mathrm{i} v)^{2}$. Thus, the curves in i), ii), and iii) are satisfying (18).

Corollary 3. Let $x: I \rightarrow \mathbb{Q}^{5} \subset \mathbb{E}_{1}^{6}$ be a curve with constant cone curvature in the lightlike cone $\mathbb{Q}^{5}$. Assume

$$
\kappa_{4} \tau_{2}+\kappa_{2} \tau_{3}=0
$$

## Consider the equation

$$
\begin{align*}
& \quad t^{2}+\left(\tau_{1}^{2}+\tau_{2}^{2}+\tau_{3}^{2}-2 \kappa_{1}\right) t \\
& -\left(\kappa_{2}^{2}+\kappa_{4}^{2}+\left(2 \kappa_{1}-\tau_{1}^{2}\right)\left(\tau_{2}^{2}+\tau_{3}^{2}\right)+\left(\kappa_{3}+\tau_{1} \tau_{2}\right)^{2}\right)=0 \tag{24}
\end{align*}
$$

i) If the roots of (24) are $\mu>0$ and $v<0$, then

$$
\begin{aligned}
x(s)= & a_{1} \sinh (s \sqrt{\mu})+a_{2} \cosh (s \sqrt{\mu})+a_{3} \sin (s \sqrt{-v}) \\
& +a_{4} \cos (\sqrt{-v} s)+a_{5} s+a_{6} .
\end{aligned}
$$

ii) If at least one of the roots of (24) is zero, then

$$
\begin{equation*}
x(s)=a_{1} s^{2}+a_{2} s+a_{3} \tag{25}
\end{equation*}
$$

provided $\lambda=0$,

$$
\begin{equation*}
x(s)=a_{1} \sinh (s \sqrt{2 \lambda})+a_{2} \cosh (s \sqrt{2 \lambda})+a_{3} \tag{26}
\end{equation*}
$$

provided $\lambda>0$, and

$$
\begin{equation*}
x(s)=a_{1} \sin (s \sqrt{-2 \lambda})+a_{2} \cos (s \sqrt{-2 \lambda})+a_{3} \tag{27}
\end{equation*}
$$

provided $\lambda<0$, where $\lambda:=2 \kappa_{1}-\tau_{1}^{2}$.
Here, $a_{i} \in \mathbb{E}_{1}^{6}, i \in\{1,2,3,4,5,6\}$, are suitable constant vectors.

Proof. The discriminant of (24) is

$$
\begin{aligned}
\Delta=\left(\tau_{2}^{2}+\tau_{3}^{2}+2 \kappa_{1}-\right. & \left.\tau_{1}^{2}\right)^{2} \\
& +4 \kappa_{2}^{2}+4 \kappa_{4}^{2}+4\left(\kappa_{3}+\tau_{1} \tau_{2}\right)^{2} \geq 0
\end{aligned}
$$

First, assume $\Delta>0$. If the roots of (24) are $\mu>0$ and $v<0$, then, as in Corollary 1, we may prove i). If one of the roots of (24) is zero, then (24) yields

$$
\begin{gathered}
\kappa_{2}^{2}+\kappa_{4}^{2}+\left(2 \kappa_{1}-\tau_{1}^{2}\right)\left(\tau_{2}^{2}+\tau_{3}^{2}\right)+\left(\kappa_{3}+\tau_{1} \tau_{2}\right)^{2}=0, \\
\kappa_{2}=\kappa_{4}=\kappa_{3}+\tau_{1} \tau_{2}=\left(2 \kappa_{1}-\tau_{1}^{2}\right)\left(\tau_{2}^{2}+\tau_{3}^{2}\right)=0,
\end{gathered}
$$

from which together with (20), we have

$$
x^{\prime \prime \prime}=\left(2 \kappa_{1}-\tau_{1}^{2}\right) x^{\prime} .
$$

The solution of this equation in the different cases is (25), (26), and (27). In the case $\Delta=0$, we conclude that any two roots of (24) are zero, and the curve is then (25).

Theorem 4. Let $x: I \rightarrow \mathbb{Q}^{6} \subset \mathbb{E}_{1}^{7}$ be a curve in the lightlike cone $\mathbb{Q}^{6}$. If the cone curvature functions of the curve are constant, then the curve satisfies the differential equation

$$
\begin{equation*}
x^{(7)}+A x^{(5)}+B x^{(3)}+C x^{\prime}=0, \tag{28}
\end{equation*}
$$

where $A, B$, and $C$ are constant coefficients.
Proof. Proceeding as in the proofs of Theorems 2 and 3, we obtain the other derivatives of $x$ as

$$
\begin{align*}
x^{\prime \prime \prime}= & \tau_{1} \kappa_{2} x+\left(2 \kappa_{1}-\tau_{1}^{2}\right) x^{\prime}+\kappa_{2} \alpha_{2} \\
& +\left(\kappa_{3}+\tau_{1} \tau_{2}\right) \alpha_{3}+\kappa_{4} \alpha_{4}+\kappa_{5} \alpha_{5},  \tag{29}\\
x^{(4)}= & \left(\kappa_{2}^{2}+\kappa_{3}^{2}+\kappa_{4}^{2}+\kappa_{5}^{2}+\kappa_{3} \tau_{1} \tau_{2}\right) x+\left(2 \kappa_{1}-\tau_{1}^{2}\right) x^{\prime \prime} \\
& -\left(\kappa_{3} \tau_{2}+\tau_{1} \tau_{2}^{2}\right) \alpha_{2}+\left(\kappa_{2} \tau_{2}-\kappa_{4} \tau_{3}\right) \alpha_{3} \\
& +\left(\kappa_{3} \tau_{3}+\tau_{1} \tau_{2} \tau_{3}-\kappa_{5} \tau_{4}\right) \alpha_{4}+\kappa_{4} \tau_{4} \alpha_{5}, \\
x^{(5)}= & \left(\kappa_{4} \tau_{1} \tau_{2} \tau_{3}-\kappa_{2} \tau_{1} \tau_{2}^{2}\right) x \\
& +\left(\kappa_{2}^{2}+\kappa_{3}^{2}+\kappa_{4}^{2}+\kappa_{5}^{2}+\tau_{1}^{2} \tau_{2}^{2}+2 \kappa_{3} \tau_{1} \tau_{2}\right) x^{\prime} \\
& +\left(2 \kappa_{1}-\tau_{1}^{2}\right) x^{\prime \prime \prime}+\left(\kappa_{4} \tau_{2} \tau_{3}-\kappa_{2} \tau_{2}^{2}\right) \alpha_{2} \\
& -\left(\left(\kappa_{3}+\tau_{1} \tau_{2}\right)\left(\tau_{2}^{2}+\tau_{3}^{2}\right)-\kappa_{5} \tau_{3} \tau_{4}\right) \alpha_{3} \\
& +\left(\kappa_{2} \tau_{2} \tau_{3}-\kappa_{4} \tau_{3}^{2}-\kappa_{4} \tau_{4}^{2}\right) \alpha_{4} \\
& +\left(\kappa_{3} \tau_{3} \tau_{4}+\tau_{1} \tau_{2} \tau_{3} \tau_{4}-\kappa_{5} \tau_{4}^{2}\right) \alpha_{5},  \tag{30}\\
x^{(6)}= & \left(2 \kappa_{2} \kappa_{4} \tau_{2} \tau_{3}+2 \kappa_{3} \kappa_{5} \tau_{3} \tau_{4}-\tau_{2}^{2}\left(\kappa_{2}^{2}+\kappa_{3}^{2}+\kappa_{3} \tau_{1} \tau_{2}\right)\right. \\
& -\tau_{3}^{2}\left(\kappa_{3}^{2}+\kappa_{4}^{2}+\kappa_{3} \tau_{1} \tau_{2}\right)-\tau_{4}^{2}\left(\kappa_{4}^{2}+\kappa_{5}^{2}\right)
\end{align*}
$$

$$
\begin{align*}
&\left.+\kappa_{5} \tau_{1} \tau_{2} \tau_{3} \tau_{4}\right) x \\
&+\left(\kappa_{2}^{2}+\kappa_{3}^{2}+\kappa_{4}^{2}+\kappa_{5}^{2}+\tau_{1}^{2} \tau_{2}^{2}+2 \kappa_{3} \tau_{1} \tau_{2}\right) x^{\prime \prime} \\
&+\left(2 \kappa_{1}-\tau_{1}^{2}\right) x^{(4)} \\
&+\left(\left(\tau_{2}^{2}+\tau_{3}^{2}\right)\left(\kappa_{3} \tau_{2}+\tau_{1} \tau_{2}^{2}\right)-\kappa_{5} \tau_{2} \tau_{3} \tau_{4}\right) \alpha_{2} \\
&+\left.\left(\tau_{2}^{2}+\tau_{3}^{2}\right)\left(\kappa_{4} \tau_{3}-\kappa_{2} \tau_{2}\right)+\kappa_{4} \tau_{3} \tau_{4}^{2}\right) \alpha_{3} \\
&-\left(\tau_{2}^{2}+\tau_{3}^{2}\right)\left(\kappa_{3} \tau_{3}+\tau_{1} \tau_{2} \tau_{3}\right) \\
&\left.+\tau_{4}^{2}\left(\kappa_{3} \tau_{3}-\kappa_{5} \tau_{4}+\tau_{1} \tau_{2} \tau_{3}\right)\right) \alpha_{4} \\
&+\left(\kappa_{2} \tau_{2} \tau_{3} \tau_{4}-\kappa_{4} \tau_{3}^{2} \tau_{4}-\kappa_{4} \tau_{4}^{3}\right) \alpha_{5} \\
& x^{(7)}=\left(\left(\tau_{2}^{2}+\tau_{3}^{2}+\tau_{4}^{2}\right)\left(\kappa_{2} \tau_{1} \tau_{2}^{2}-\kappa_{4} \tau_{1} \tau_{2} \tau_{3}\right)-\kappa_{4} \tau_{1} \tau_{2}^{2} \tau_{4}^{2}\right) x \\
&+\left(2 \kappa_{2} \kappa_{4} \tau_{2} \tau_{3}+2 \kappa_{3} \kappa_{5} \tau_{3} \tau_{4}-2 \kappa_{3} \tau_{1} \tau_{2}\left(\tau_{2}^{2}+\tau_{3}^{2}\right)\right. \\
&-\kappa_{2}^{2} \tau_{2}^{2}-\kappa_{3}^{2}\left(\tau_{2}^{2}+\tau_{3}^{2}\right)-\kappa_{4}^{2}\left(\tau_{3}^{2}+\tau_{4}^{2}\right) \\
&+\kappa_{5} \tau_{1} \tau_{2} \tau_{3} \tau_{4}-\kappa_{5}^{2} \tau_{4}^{2} \\
&\left.-\tau_{1}^{2} \tau_{2}^{2}\left(\tau_{2}^{2}+\tau_{3}^{2}\right)+\kappa_{5} \tau_{1}^{2} \tau_{3} \tau_{4}\right) x^{\prime} \\
&+\left(\kappa_{2}^{2}+\kappa_{3}^{2}+\kappa_{4}^{2}+\kappa_{5}^{2}+\tau_{1}^{2} \tau_{2}^{2}\right. \\
&\left.-2 \kappa_{3} \tau_{1} \tau_{2}\left(\tau_{2}^{2}+\tau_{3}^{2}\right)+2 \kappa_{3} \tau_{1} \tau_{2}\right) x^{\prime \prime \prime} \\
&+\left(2 \kappa_{1}-\tau_{1}^{2}\right) x^{(5)} \\
&+\left(\left(\tau_{2}^{2}+\tau_{3}^{2}+\tau_{4}^{2}\right)\left(\kappa_{2} \tau_{2}^{2}-\kappa_{4} \tau_{2} \tau_{3}\right)-\kappa_{2} \tau_{2}^{2} \tau_{4}^{2}\right) \alpha_{2} \\
&+\left(\left(\tau_{2}^{2}+\tau_{3}^{2}\right)\left(\kappa_{3} \tau_{2}^{2}+\tau_{1} \tau_{2}^{3}\right)-\kappa_{5} \tau_{3} \tau_{4}\left(\left(\tau_{2}^{2}+\tau_{3}^{2}+\tau_{4}^{2}\right)\right.\right. \\
&\left.+\left(\kappa_{3}+\tau_{1} \tau_{2}\right) \tau_{3}^{2} \tau_{4}^{2}\right) \alpha_{3} \\
&+\left(\left(\kappa_{4} \tau_{3}^{2}+\kappa_{4} \tau_{4}^{2}-\kappa_{2} \tau_{2} \tau_{3}\right)\left(\tau_{2}^{2}+\tau_{3}^{2}+\tau_{4}^{2}\right)\right. \\
&\left.-\kappa_{4} \tau_{2}^{2} \tau_{4}^{2}\right){\alpha_{4}}_{-}\left(\left(\tau_{2}^{2}+\tau_{3}^{2}+\tau_{4}^{2}\right)\left(\kappa_{3} \tau_{3} \tau_{4}+\tau_{1} \tau_{2} \tau_{3} \tau_{4}-\kappa_{5} \tau_{4}^{2}\right)\right. \\
&+\left.\kappa_{5} \tau_{2}^{2} \tau_{4}^{2}\right){\alpha_{5}}
\end{align*}
$$

If we multiply (30) by $\left(\tau_{2}^{2}+\tau_{3}^{2}+\tau_{4}^{2}\right)$, multiply (29) by $\tau_{2}^{2} \tau_{4}^{2}$, and then add the resulting equations to (31), then we obtain (28), where $A, B$, and $C$ are suitable constant coefficients.

Theorem 5. Let $x: I \rightarrow \mathbb{Q}^{n+1}$ be a curve with constant cone curvature in $\mathbb{Q}^{n+1}$. Assume that the lightlike vector field $y(s)$ is as in (2).
i) If $\kappa_{i}=0$ for all $i \in\{2, \ldots, n\}$, then

$$
x(s)=a_{1} s^{2}+a_{2} s+a_{3}
$$

provided $\kappa_{1}=0$,

$$
x(s)=a_{1} \sinh \left(s \sqrt{2 \kappa_{1}}\right)+a_{2} \cosh \left(s \sqrt{2 \kappa_{1}}\right)+a_{3}
$$

provided $\kappa_{1}>0$, and

$$
x(s)=a_{1} \sin \left(s \sqrt{-2 \kappa_{1}}\right)+a_{2} \cos \left(s \sqrt{-2 \kappa_{1}}\right)+a_{3}
$$

provided $\kappa_{1}<0$.
ii) If one of the $\kappa_{i} \neq 0$ for $i \in\{2, \ldots, n\}$, then

$$
\begin{align*}
x(s)= & a_{1} \sinh (\mu s)+a_{2} \cosh (\mu s)  \tag{32}\\
& +a_{3} \sin (v s)+a_{4} \cos (v s),
\end{align*}
$$

where $a_{i} \in \mathbb{E}_{1}^{n+2}, i \in\{1,2,3\}$, are suitable vectors and $\pm \mu$ and $\pm i v$ are the real and imaginary roots of the equation

$$
t^{4}-2 \kappa_{1} t^{2}-\sum_{i=2}^{n} \kappa_{i}^{2}=0
$$

Proof. First, we prove i). By Proposition 1, all of the $\tau_{i}=0$, and because of $\kappa_{i}=0$ for all $i \in\{2, \ldots, n\}$, by the Frenet formulas (1), we conclude that

$$
x^{\prime \prime \prime}(s)-2 \kappa_{1} x^{\prime}(s)=0
$$

which has its solutions in the stated forms. Finally, we prove ii). By Proposition 1, all of the $\tau_{i}=0$, thus, by the Frenet formulas (1), we conclude that

$$
x^{(4)}(s)-2 \kappa_{1} x^{\prime \prime}(s)-\sum_{i=2}^{n} \kappa_{i}^{2} x(s)=0
$$

which has its solutions in the form (32).

## 5 Some Results for Spacelike Curves in $\mathbb{Q}^{3}$

In this section, we follow the relationship between Frenet curvatures and cone curvature functions on $\mathbb{Q}^{3}$ in special cases. J. Walrave in [9] has classified all spacelike curves in $\mathbb{E}_{1}^{4}$, and we consider such constant cone curvature curves. By Corollary 1, the only nonstraightline curves in $\mathbb{Q}^{n+1}$ are spacelike curves, thus we will not study the other types.
M. P. Torgašev and E. Šućurović in [8, Remarks 3.2 and 3.3] have proved that if $x: I \rightarrow \mathbb{Q}^{3} \subset \mathbb{E}_{1}^{4}$ is a curve in the lightlike cone $\mathbb{Q}^{3}$ with constant Frenet curvatures and the principal vector $N$ or the binormal vector $B_{1}$ is timelike, then $\bar{\tau}^{2}=\bar{\sigma}^{2}$. In the following theorems, in the cases when $B_{2}$ or $N$ is timelike, we independently prove this, and next we will show that if $B_{1}$ is timelike or $B_{1}, B_{2}$ are lightlike, then $\kappa_{2}=\bar{\tau}=\bar{\sigma}=0$.

Remark. In the sequel, we will write $\kappa_{1}, \kappa_{2}$, and $\tau$ for cone curvature functions, and $\bar{\kappa}, \bar{\tau}$, and $\bar{\sigma}$ for Frenet curvatures functions.

Theorem 6. Let $x: I \rightarrow \mathbb{Q}^{3} \subset \mathbb{E}_{1}^{4}$ be a curve in the lightlike cone $\mathbb{Q}^{3}$. If the cone curvature functions of the curve are constant, then its Frenet curvatures are constant.

Proof. Case 1. Let $N$ and $B_{1}$ be spacelike. Thus, $B_{2}$ is timelike. In this case, we have the Frenet formulas

$$
\left(\begin{array}{c}
T  \tag{33}\\
N \\
B_{1} \\
B_{2}
\end{array}\right)^{\prime}=\left(\begin{array}{cccc}
0 & \bar{\kappa} & 0 & 0 \\
-\bar{\kappa} & 0 & \bar{\tau} & 0 \\
0 & -\bar{\tau} & 0 & \bar{\sigma} \\
0 & 0 & \bar{\sigma} & 0
\end{array}\right)\left(\begin{array}{c}
T \\
N \\
B_{1} \\
B_{2}
\end{array}\right)
$$

By (1), for $\mathbb{Q}^{3}$, we have

$$
\begin{gather*}
N=\frac{1}{\bar{\kappa}} \alpha_{1}^{\prime}=\frac{1}{\bar{\kappa}}\left(\kappa_{1} x-y+\tau \alpha_{2}\right)  \tag{34}\\
\bar{\kappa}^{2}=\tau^{2}-2 \kappa_{1}
\end{gather*}
$$

Thus, the first Frenet curvature of the curve is constant. By (33) and (34), we have

$$
\begin{gather*}
N^{\prime}=-\bar{\kappa} T+\bar{\tau} B_{1}=-\bar{\kappa} T+\frac{\tau \kappa_{2}}{\bar{\kappa}} x+\frac{\kappa_{2}}{\bar{\kappa}} \alpha_{2}  \tag{35}\\
\bar{\tau}^{2}=\left(\frac{\kappa_{2}}{\bar{\kappa}}\right)^{2}, B_{1}=\varepsilon\left(\tau x+\alpha_{2}\right)
\end{gather*}
$$

where $\varepsilon= \pm 1$. Thus, $\bar{\tau}$ is constant. Now by (35) and the third equation of (33), we have

$$
B_{1}^{\prime}=-\bar{\tau} N+\bar{\sigma} B_{2}=\varepsilon \kappa_{2} x,
$$

so that $B_{1}^{\prime}$ is lightlike, and as a result of orthonormality, we have

$$
\bar{\tau}^{2}-\bar{\sigma}^{2}=0, \bar{\tau}^{2}=\bar{\sigma}^{2}
$$

and hence $\bar{\sigma}$ is constant.
Case 2. Let $B_{1}, B_{2}$ be spacelike and $N$ be timelike. In this case, we have the Frenet formulas

$$
\left(\begin{array}{c}
T \\
N \\
B_{1} \\
B_{2}
\end{array}\right)^{\prime}=\left(\begin{array}{cccc}
0 & \bar{\kappa} & 0 & 0 \\
\bar{\kappa} & 0 & \bar{\tau} & 0 \\
0 & \bar{\tau} & 0 & \bar{\sigma} \\
0 & 0 & -\bar{\sigma} & 0
\end{array}\right)\left(\begin{array}{c}
T \\
N \\
B_{1} \\
B_{2}
\end{array}\right) .
$$

Similarly as in Case 1, this results in the Frenet curvatures being constant. In this case, we have $\bar{\kappa}^{2}=2 \kappa_{1}-\tau^{2}$.

Case 3. Let $N, B_{2}$ be spacelike and $B_{1}$ be timelike. We have the Frenet equations

$$
\left(\begin{array}{c}
T  \tag{36}\\
N \\
B_{1} \\
B_{2}
\end{array}\right)^{\prime}=\left(\begin{array}{cccc}
0 & \bar{\kappa} & 0 & 0 \\
-\bar{\kappa} & 0 & \bar{\tau} & 0 \\
0 & \bar{\tau} & 0 & \bar{\sigma} \\
0 & 0 & \bar{\sigma} & 0
\end{array}\right)\left(\begin{array}{c}
T \\
N \\
B_{1} \\
B_{2}
\end{array}\right) .
$$

This case is similar to Cases 1 and 2 as well, and by an elementary calculation, we conclude that $\bar{\kappa}^{2}=\tau^{2}-2 \kappa_{1}$ and $\kappa_{2}=\bar{\tau}=\bar{\sigma}=0$ are constant (see [8, Remark 3.2 and Theorem 3.7]).

Case 4 . Let $N$ be spacelike and $B_{1}, B_{2}$ be lightlike. In this case, we have

$$
B_{1}=N^{\prime}+\bar{\tau} T,\left\langle B_{1}, B_{2}\right\rangle=1,\left\langle B_{1}, B_{1}\right\rangle=\left\langle B_{2}, B_{2}\right\rangle=0
$$

and the Frenet formulas are

$$
\left(\begin{array}{c}
T \\
N \\
B_{1} \\
B_{2}
\end{array}\right)^{\prime}=\left(\begin{array}{cccc}
0 & \bar{\kappa} & 0 & 0 \\
-\bar{\kappa} & 0 & \bar{\tau} & 0 \\
0 & 0 & \bar{\sigma} & 0 \\
0 & -\bar{\tau} & 0 & -\bar{\sigma}
\end{array}\right)\left(\begin{array}{c}
T \\
N \\
B_{1} \\
B_{2}
\end{array}\right) .
$$

This case is similar to Case 3 . We conclude that $\bar{\kappa}^{2}=\tau^{2}-$ $2 \kappa_{1}$ and $\kappa_{2}=\bar{\tau}=\bar{\sigma}=0$ are constant.

Corollary 4. Let $x: I \rightarrow \mathbb{Q}^{3} \subset \mathbb{E}_{1}^{4}$ be the curve in Cases 3 or 4 of Theorem 6. If its Frenet curvatures are constant, then it is a planar curve.

Proof. By Theorem 6, in Cases 3 or 4, we have $\kappa_{2}=0$. Thus, by [8, Theorem 3.1], the curve is planar.

Theorem 7. Let $x: I \rightarrow \mathbb{Q}^{3} \subset \mathbb{E}_{1}^{4}$ be a curve in the lightlike cone $\mathbb{Q}^{3}$ with constant Frenet curvatures. Then its cone curvature functions satisfy the relations

$$
\bar{\kappa}^{2}=\tau^{2}-2 \kappa_{1}, \kappa_{2}+\tau^{\prime}= \pm \bar{\tau} \bar{\kappa}
$$

Moreover, in Cases 3 and 4 of Theorem 6, if $\bar{\kappa} \neq 0$, then $\bar{\tau}=\bar{\sigma}=0$.

Proof. In all of the cases, by the second equation in (1), we have

$$
\alpha_{1}^{\prime}=\kappa_{1} x-y+\tau \alpha_{2}, \bar{\kappa}^{2}=\tau^{2}-2 \kappa_{1} .
$$

In Case 1, by the Frenet equation (33), we have
$N^{\prime}=-\bar{\kappa} T+\bar{\tau} B_{1}=\frac{1}{\bar{\kappa}}\left(-\bar{\kappa}^{2} x^{\prime}+\kappa_{1}^{\prime} x+\tau^{\prime} \alpha_{2}+\tau \kappa_{2} x+\kappa_{2} \alpha_{2}\right)$.
Since $B_{1}$ is spacelike, we get

$$
(\overline{\tau \bar{\kappa}})^{2}=\left(\kappa_{2}+\tau^{\prime}\right)^{2}
$$

Again, by the Frenet equation (33), we have

$$
\begin{aligned}
B_{1}^{\prime}= & -\bar{\tau} N+\bar{\sigma} B_{2} \\
= & \frac{1}{\bar{\tau} \kappa}\left(\left(\kappa_{2}^{2}+2 \tau^{\prime} \kappa_{2}+\tau \kappa_{2}^{\prime}+\kappa_{1}^{\prime \prime}\right) x\right. \\
& \left.+\left(\tau \kappa_{2}+\kappa_{1}^{\prime}-\tau \kappa_{2}-\tau \tau^{\prime}\right) x^{\prime}+\left(\kappa_{2}^{\prime}+\tau^{\prime \prime}\right) \alpha_{2}\right\}
\end{aligned}
$$

and since $\tau \tau^{\prime}-\kappa_{1}^{\prime}=\kappa_{2}^{\prime}+\tau^{\prime \prime}=0$ and $\kappa_{1}^{\prime \prime}=\tau \tau^{\prime \prime}+\tau^{\prime 2}$, we conclude that

$$
-\bar{\tau} N+\bar{\sigma} B_{2}=(\overline{\tau \kappa}) x .
$$

Since $x$ is lightlike, $N$ is spacelike and $B_{2}$ is timelike, thus $\bar{\tau}^{2}=\bar{\sigma}^{2}$.

Case 2 is similar.
In Case 3, by the Frenet equation (36), we have
$N^{\prime}=-\bar{\kappa} T+\bar{\tau} B_{1}=\frac{1}{\bar{\kappa}}\left(-\bar{\kappa} x^{\prime}+\kappa_{1}^{\prime} x+\tau^{\prime} \alpha_{2}+\tau \kappa_{2} x+\kappa_{2} \alpha_{2}\right)$.
Since $B_{1}$ is timelike, we have

$$
-(\overline{\tau \kappa})^{2}=\left(\kappa_{2}+\tau^{\prime}\right)^{2} .
$$

Since $\bar{\kappa} \neq 0$, we obtain $\bar{\tau}=0$.
Corollary 5. Let $x: I \rightarrow \mathbb{Q}^{3} \subset \mathbb{E}_{1}^{4}$ be a curve in the lightlike cone $\mathbb{Q}^{3}$. Assume that the lightlike vector field y is defined in (2). Then the cone curvature functions of the curve are constant if and only if its Frenet curvatures are constant.

Proof. In this case, from $\langle N, N\rangle=1$, we have

$$
\bar{\kappa}^{2}=-2 \kappa_{1} .
$$

Thus the first cone curvature function is constant if and only if the first Frenet curvature function is constant. In Case 1 of Theorem 6, we conclude

$$
\bar{\tau}^{2}=\left(\frac{\kappa_{2}}{\bar{\kappa}}\right)^{2}
$$

Hence the second cone curvature function is constant if and only if the second (third) Frenet curvature is constant.
H. Liu in [4] proved that if $x: I \rightarrow \mathbb{Q}^{2}$ is a helix such that its velocity vector field has constant angle with a constant vector, then its cone curvature function satisfies

$$
\kappa(s)=c_{1}\left(s+c_{2}\right)^{-2}
$$

where $c_{1} \neq 0$ and $c_{2}$ are constant. Now, we prove a similar result as follows.

Theorem 8. Let $x: I \rightarrow \mathbb{Q}^{3}$ be a curve such that $\alpha_{1}$ and $\alpha_{2}$ have constant angle with the constant vector $b$. Then its cone curvature function $\kappa_{1}$ and $\kappa_{2}$ satisfy

$$
\kappa_{1} l+\kappa_{2} \tilde{l}=c_{1}\left(s+c_{2}\right)^{-2}
$$

where $\left\langle\alpha_{1}, b\right\rangle=l,\left\langle\alpha_{2}, b\right\rangle=\tilde{l}$, and $c_{1}, c_{2}$ are real constants.
Proof. From $\left\langle\alpha_{1}, b\right\rangle=l$, we have

$$
\langle x, b\rangle=l s+l_{0}
$$

and

$$
0=\left\langle\alpha_{1}^{\prime}, b\right\rangle=\left\langle\kappa_{1} x-y-\tau \alpha_{2}, b\right\rangle=\kappa_{1}\langle x, b\rangle-\langle y, b\rangle+\tau \tilde{l} .
$$

Thus,

$$
\begin{equation*}
\kappa_{1}\left(l s+l_{0}\right)+2 \kappa_{1} l+\left(\kappa_{2}+\tau^{\prime}\right) \tilde{l}=0 \tag{37}
\end{equation*}
$$

Similarly, $\left\langle\alpha_{2}, b\right\rangle=\tilde{l}$ yields

$$
\left\langle\alpha_{2}^{\prime}, b\right\rangle=\kappa_{2}\left(l s+l_{0}\right)-\tau l=0
$$

By differentiation, we have

$$
\begin{equation*}
\tau^{\prime}=\kappa_{2}^{\prime} s+\kappa_{2}+\kappa_{2}^{\prime} \frac{l_{0}}{l} \tag{38}
\end{equation*}
$$

From (38) and (37), we get

$$
\left(\kappa_{1}^{\prime} l+\kappa_{2}^{\prime} \tilde{l}\right)\left(l s+l_{0}\right)+2\left(\kappa_{1} l+\kappa_{2}^{\prime} \tilde{l}\right)=0
$$

so that

$$
\frac{\kappa_{1}^{\prime} l+\kappa_{2}^{\prime} \tilde{l}}{\kappa_{1} l+\kappa_{2} \tilde{l}}=-2 \frac{l}{l s+l_{0}}
$$

Hence the solutions of this differential equation are

$$
\kappa_{1} l+\kappa_{2} \tilde{l}=c\left(l s+l_{0}\right)^{-2} .
$$

Thus,

$$
\kappa_{1} l+\kappa_{2} \tilde{l}=c_{1}\left(s+c_{2}\right)^{-2}
$$

completing the proof.

Theorem 9. Let $x: I \rightarrow \mathbb{Q}^{3}$ be a spacelike curve such that the position vector $x(s)$ has constant inner product with a constant vector $b$. If the lightlike vector field $y(s)$ is defined as in (2), then the cone curvature functions $\kappa_{1}$ and $\kappa_{2}$ satisfy

$$
\begin{equation*}
\kappa_{2}=-\left(\frac{\kappa_{1}^{\prime}}{\kappa_{2}}\right)^{\prime} \tag{39}
\end{equation*}
$$

Proof. Let $\langle x, b\rangle=l$ such that $l$ is a constant. Then

$$
\left\langle x^{\prime}, b\right\rangle=0
$$

By the Frenet equations (1), we have

$$
\langle y, b\rangle=\kappa_{1} l
$$

The vector $b$, by the frame $\left\{x, x^{\prime}, y, \alpha_{2}\right\}$, can be expressed as

$$
b=l \kappa_{1}(s) x(s)+l y(s)+\lambda(s) \alpha_{2}(s)
$$

Thus, $b^{\prime}=0$ and

$$
\left(l \kappa_{1}^{\prime}(s)+\lambda(s) \kappa_{2}(s)\right) x(s)+\left(\lambda^{\prime}(s)-l \kappa_{2}(s)\right) \alpha_{2}(s)=0
$$

Therefore,

$$
\lambda^{\prime}(s)-l \kappa_{2}(s)=0, l \kappa_{1}^{\prime}(s)+\lambda(s) \kappa_{2}(s)=0
$$

which yields (39).

## Acknowledgement

The first and last authors thank the University of Mohaghegh Ardabili for supporting this research.

## References

[1] Nemat Abazari, Martin Bohner, Ilgın Sağer, and Yusuf Yayli, Stationary acceleration of Frenet curves, J. Inequal. Appl., Paper No. 92, 13 pages, 2017.
[2] Mehmet Bektaş and Mihriban Külahci, Differential equations characterizing spacelike curves in the 3dimensional lightlike cone, Palest. J. Math., Vol. 6, No. 2, pp. 330-337 (2017).
[3] Wolfgang Kühnel, Differential geometry, volume 77 of Student Mathematical Library. American Mathematical Society, Providence, RI, 2015. Curves-surfacesmanifolds, Third edition [of MR1882174], Translated from the 2013 German edition by Bruce Hunt, with corrections and additions by the author.
[4] Huili Liu, Curves in the lightlike cone, Beiträge Algebra Geom., Vol. 45, No. 1, pp. 291-303 (2004).
[5] Huili Liu and Qingxian Meng, Representation formulas of curves in a two- and three-dimensional lightlike cone, Results Math., Vol. 59, No. 3-4, pp. 437-451 (2011).
[6] Rafael López, Differential geometry of curves and surfaces in Lorentz-Minkowski space, Int. Electron. J. Geom., Vol. 7, No. 1, pp. 44-107 (2014).
[7] Barrett O'Neill, Semi-Riemannian geometry, volume 103 of Pure and Applied Mathematics. Academic Press, Inc. [Harcourt Brace Jovanovich, Publishers], New York, 1983. With applications to relativity.
[8] Miroslava Petrović-Torgašev and Emilija Šućurović, Wcurves in Minkowski space-time, Novi Sad J. Math., Vol. 32, No. 2, pp. 55-65 (2002).
[9] Johan Walrave, Curves and surfaces in Minkowski space, ProQuest LLC, Ann Arbor, MI, 1995. Thesis (Ph.D.)Katholieke Universiteit Leuven (Belgium).
[10] Suha Yilmaz and Melih Turgut, On the differential geometry of the curves in Minkowski space-time. I, Int. J. Contemp. Math. Sci., Vol. 3, No. 25-28, pp. 1343-1349 (2008).


Nemat Abazari was born in Ardabil, Iran, in 1972. He received the BS degree from the University of Tabriz, Iran, in 1994, the MS degree from the Valiasr University of Rafsanjan, Rafsanjan, Iran, 2001, and the PhD degree in Geometry from Ankara University, Ankara, Turkey, in 2011. Since 2011, he is Associate Professor at the University of Mohaghegh Ardabili, Ardabil, Iran.


Martin Bohner is the Curators' Professor of Mathematics and Statistics at Missouri University of Science and Technology in Rolla, Missouri, USA. He received the BS (1989) and MS (1993) in Econo-mathematics and PhD (1995) from Universität Ulm, Germany, and MS (1992) in Applied Mathematics from San Diego State University. His research interests center around differential, difference, and dynamic equations as well as their applications to economics, finance, biology, physics, and engineering. He is the author of six textbooks and more than 250 publications, Editor-in-Chief of four international journals, Associate Editor for more than 50 international journals, and President of ISDE, the International Society of Difference Equations. Professor Bohner's honors at Missouri S\&T include five Faculty Excellence Awards, one Faculty Research Award, and eight Teaching Awards.

Ilgın Sağer is a Teaching
 Assistant Professor at the Department of Mathematics and Computer Science at University of Missouri-St. Louis. She received the BS (2002), MS (2005), and PhD (2015) in Mathematics from Ankara University. Her research interests are symplectic geometry and the theory of elastic curves in Lorentz-Minkowski spaces.


Alireza Sedaghatdoost is a Phd student in Differential Geometry at Mohaghegh Ardebili University. He received the BS (2007) and MS (2009) in Geometry from Tabriz University. His research interests are minimal surfaces and the theory of curves in the Euclidean and Lorentz-Minkowski spaces.


[^0]:    * Corresponding author e-mail: abazari@uma.ac.ir

